

EMFCSC

International School of Subnuclear Physics

*GRAVITY AND MATTER
IN THE SUBNUCLEAR WORLD*

58th Course. ERICE 15-24 JUNE 2022

*Are Neutrinos and Antineutrinos
the same particle?*

A. Bettini

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Summary

- **Majorana vs Dirac spin $1/2$ particles**
- **Two-neutrinos and neutrino-less double beta decay**
- **Backgrounds and energy resolution**
- **Examples of leading edge experiments**
- **Limits on effective Majorana mass**
- **Hints from atomic and supramolecular physics**
- **Conclusions**

1937 *Majorana anti* $\nu=\nu$?



TEORIA SIMMETRICA DELL'ELETTRONE E DEL POSITRONE

Nota di ETTORE MAJORANA

The new approach allows to “*not only to give a symmetric form to the electron-positron theory, but also to build a substantially novel theory for the particles deprived of electric charge (... hypothetical neutrinos)*”

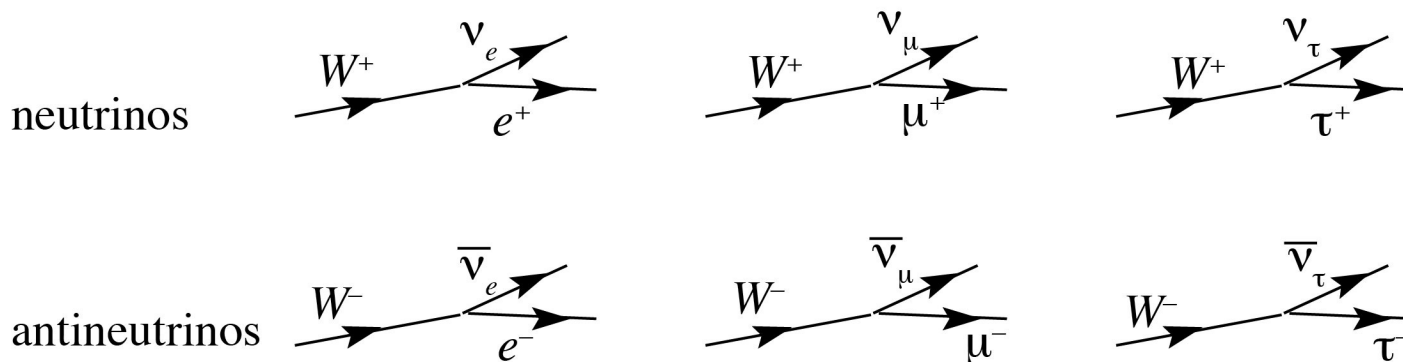
.....it is probably “*not yet possible to ask to the experience to decide between this new theory and the simple extension of the Dirac equation to the neutral particles*”

Definitions

ν_e is the neutral particle produced with an e^+ (e.g. β^+ decay)

“anti” ν_e is the neutral particle produced with an e^- (e.g. β^- decay)

Same for ν_μ and ν_τ and their “anti” particles

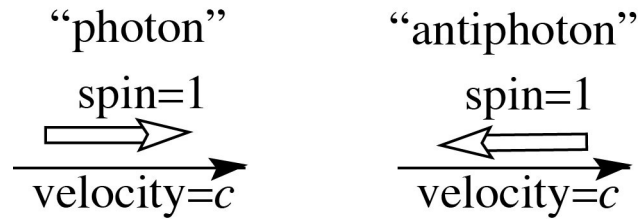


Majorana: *The advantage of this procedure....is that there is now no reason to assume the existence of ... antineutrinos.*

Matter or antimatter?

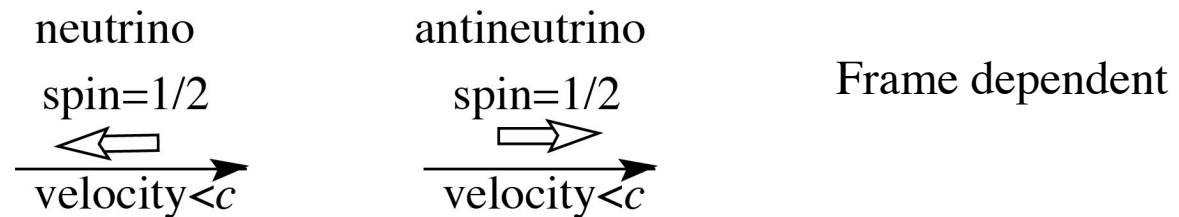
Photon is a completely neutral particle, no charge-type quantum number

The antiparticle of the photon is the photon. Just invert helicity



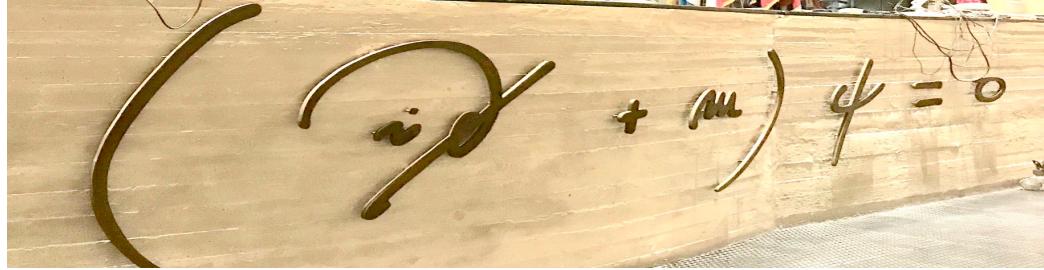
Neutrino and antineutrino in the SM are assumed to be different particles, distinguished by the lepton number. No experimental evidence

If lepton number is not a good quantum number neutrino and antineutrino may be two states of the same particle, distinguished by helicity (negative vs positive)



If neutrino and antineutrino are two states of the same particle, the matter-antimatter asymmetry in the Universe can find an explanation

Majorana vs Dirac equation



Dirac

$$\psi = \begin{pmatrix} \varphi \\ \chi \end{pmatrix}$$

$$(i\partial_t - i\vec{\sigma} \cdot \vec{\nabla})\varphi = m\chi$$

$$(i\partial_t + i\vec{\sigma} \cdot \vec{\nabla})\chi = m\varphi$$

Majorana

$$\psi_M = \begin{pmatrix} \varphi \\ i\sigma_2\varphi^* \end{pmatrix}$$

$$(i\partial_t - i\vec{\sigma} \cdot \vec{\nabla})\varphi = im\sigma_2\varphi^*$$

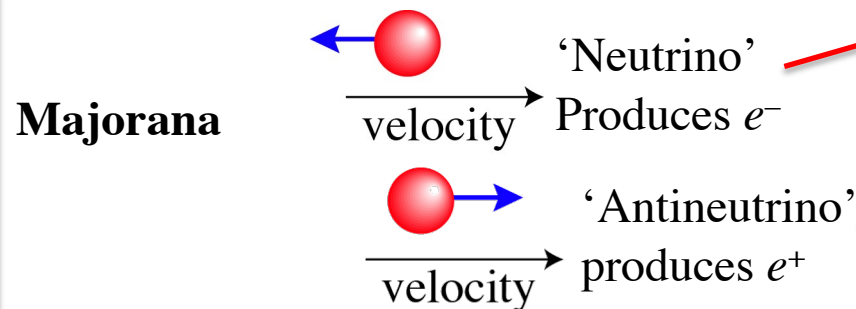
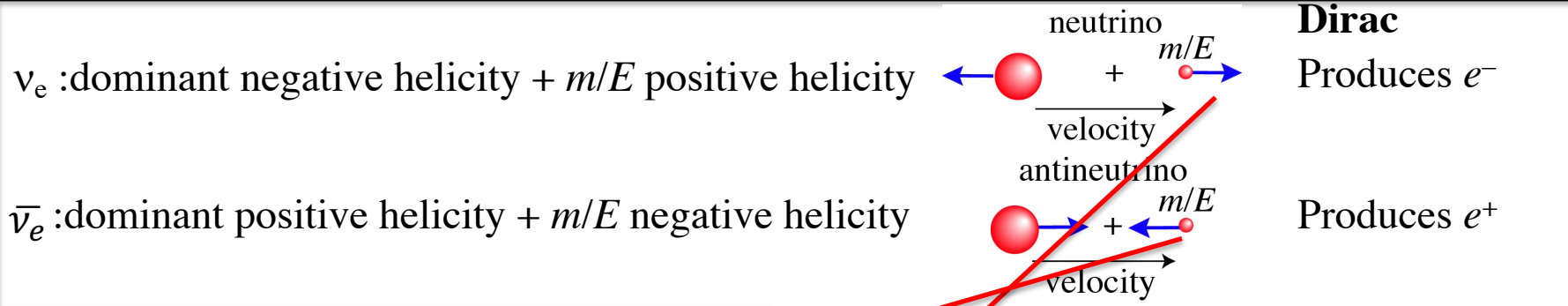
**If m=0,
Dirac =
Majorana**

$$\psi_M^C = C\gamma^0\psi_M^* = \begin{pmatrix} 0 & -i\sigma_2 \\ i\sigma_2 & 0 \end{pmatrix} \begin{pmatrix} \varphi \\ i\sigma_2\varphi^* \end{pmatrix}^* = \begin{pmatrix} \varphi \\ i\sigma_2\varphi^* \end{pmatrix} = \psi_M$$

Majorana particles are **completely neutral spin 1/2 particles** (like γ , Z, etc. for bosons)

Is neutrino completely neutral?

V-A only left-handed field ($\gamma_5 = -1$)



Only the extremely small m/E component distinguishes between Dirac and Majorana
 Beam energy $E_\nu = 10$ MeV, $m = 100$ meV
 $\Rightarrow (m/E_\nu)^2 = 10^{-14}$

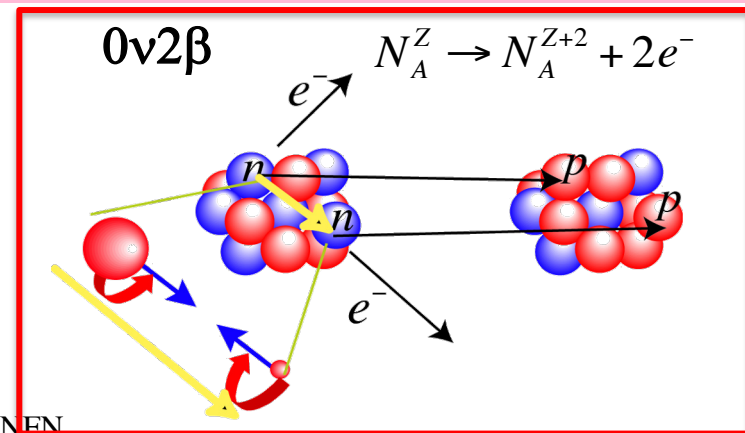
ν_e not producing e^+ do not prove lepton number conservation



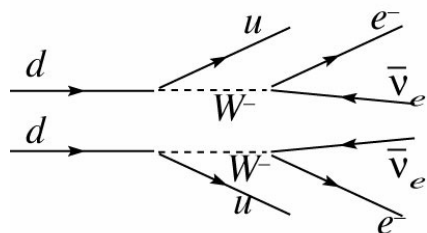
AVOGADRO, Help!

6×10^{26}

Have a large N



Double beta active isotopes



Two neutrino double beta decay
2nd order weak interaction
In nuclides stable against β decay

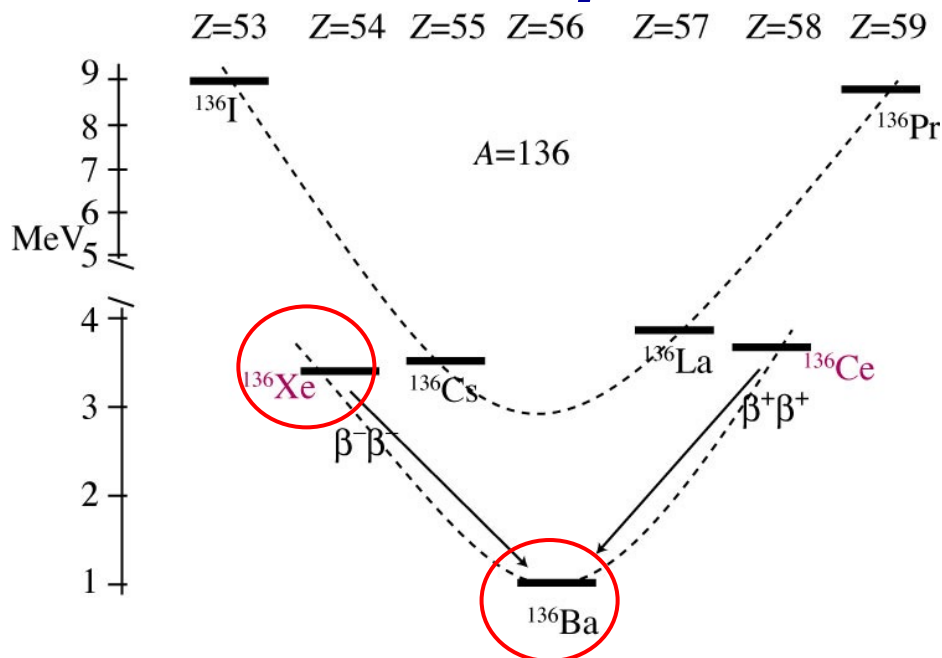
Nucleon level

$$n+n \rightarrow 2e^- + 2p + 2\bar{\nu}_e$$

Nuclear level

$$(A, Z) \rightarrow (A, Z + 2) + 2e^- + 2p + \bar{\nu}_e$$

Several have been measured



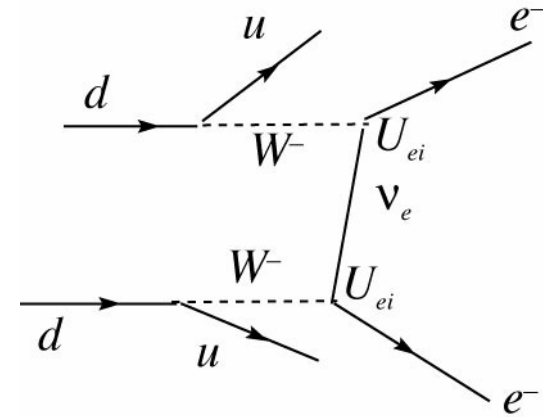
Isotope	$T^{2\nu}_{1/2}(10^{20} \text{ yr})$
⁷⁶ Ge	19.3 ± 0.9
⁸² Se	0.94 ± 0.06
¹⁰⁰ Mo	0.07 ± 0.002
¹³⁰ Te	8.2 ± 0.7
¹³⁶ Xe	21.7 ± 0.7

$2\beta 0\nu$ Decay

Majorana neutrino couples to W as Dirac neutrino

The SM violation is in the propagator

At one vertex the ($h=+$) component matters, the ($h=-$) component at the other vertex



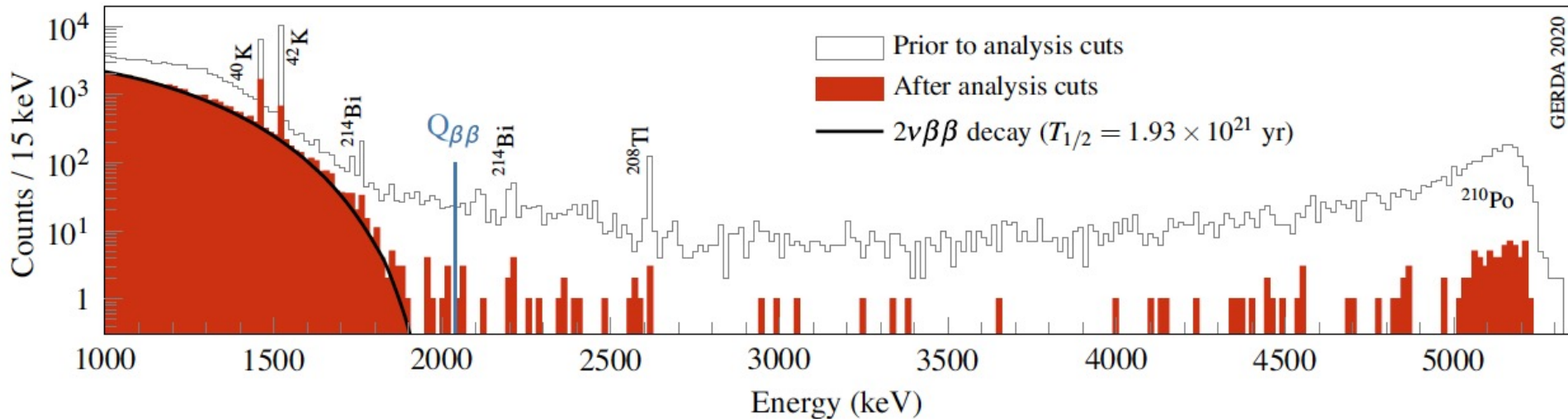
$$M_{ee} = [U_{e1}^2 m_1 + U_{e2}^2 m_2 + U_{e3}^2 m_3] \sim [0.67 m_1 + 0.30 m_2 e^{2i(\eta_1 - \eta_2)} + 0.03 m_3 e^{2i(\eta_1 + \delta_{CP})}]$$

$$1/T_{1/2} = G(Q, Z) |NME|^2 M_{ee}^2$$

Sensitivity to M_{ee} scales as the (sensitivity to $T_{1/2}$)^{1/2}

Phase-space factor, PSF, $G(Q, Z)$ calculated taking into account that electrons are in an atom
Nuclear matrix element calculation quite uncertain (factor 2-3)

The sum-energy electrons spectrum



GERDA 2020 example

On the left $2\nu 2\beta$ continuous spectrum. Potential background for $0\nu 0\beta$; need energy resolution

Signal: a line at $Q_{\beta\beta}$, which is **exactly known**; width = energy resolution ΔE (FWHM)

Backgrounds from radioactive sources reduced, by experimental cuts (see later)

Background index: b = counts per unit energy interval, unit sensitive mass, unit live time:
counts/(keV kg yr)

Figure of merit: $b\Delta E$ (counts/kg y) = total background in the region of interest (roi)

0-bkgrnd: $b\Delta E < (0.5-1)$. Sensitivity to $T_{1/2}$ grows **linearly** with exposure $M\Delta t$

non-0 bkgrnd. Sensitivity to $T_{1/2}$ grows with **sqrt** of exposure $(M\Delta t)^{1/2}$

Isotope separation

With the exception of ^{130}Te , all nuclides require isotopic enrichment

ECP. Joint Stock Company. Production Association Electrochemical Plant", Zelenogorsk, Oblast' of Oremburg (near Kransnoyarsk), Siberia, Russia; http://www.ecp.ru/index_en.shtml



Gas centrifuge processing: large number of centrifuges in series and parallel formations.

Large, background-free exposures

$0\nu 2\beta$ signal might be at $T_{1/2} = 10^{27} - 10^{28}$ yr. Large exposure needed
For a few counts of $0\nu 2\beta$ we need

- 10^{26} yr: 100 kg yr now
- 10^{27} yr: 1 t yr 2027
- 10^{28} yr: 10 t yr 2037

If the signal is a few counts the background needs to be near to 0 (< 1 count)
An order of magnitude larger sensitive mass needs an order of magnitude lower background

The low background frontier

We live in a very radioactive world
human body: $6 \cdot 10^{12}/(\text{t yr})$

Backgrounds in a FWHM per isotope ton per year.

Experiment	b ΔE ($\text{t}^{-1}\text{yr}^{-1}$)
GERDA (^{76}Ge)	1.7
EXO-200 (^{136}Xe)	139
KamLAND-ZEN (^{136}Xe)	49
NEXT-100 (^{136}Xe)	8 expected
CUORE (^{130}Te)	418

Backgrounds

1. Cosmic muons

Almost exponentially decreasing with depth

An order of magnitude for about 650 m

Short delay effects (e.g. spallation)

Detect muon and anticoincidence

Long delay effects (unstable nuclides, delayed decays)

Difficult to cut out if in the detector or in the shield

2. Neutrons from fission & (α, n) from U/Th in the rocks/environment

Depending on local geology, depth independent

Can be shielded

3. Gamma ambient flux (including radon & progeny)

Depending on local geology, depth independent

Can be shielded (but not Rn emanated inside detector)

4. Detector materials, supports, shielding, electrical connections, etc.

Use super clean components

Complete components assay

Avoid surfaces, especially in vacuum, identify surface events

Shielding example - GERDA

Water tank
moderates and absorbs neutrons

Liquid Ar cryostat
Absorbs gammas

Hyperpure copper
Absorbs radiation from steel wall

Ge detectors

Shields are made “active” detecting
backgrounds via scintillation (Ar) or
Cherenkov (water) light



G. Heusser, Ann. Rev. Nucl. Part. Sci. 45 (1995) 543.

17-Jun-22

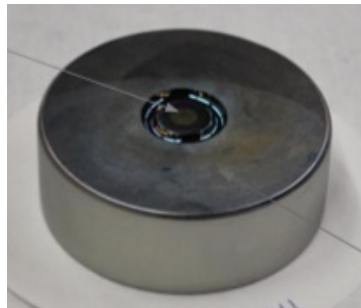
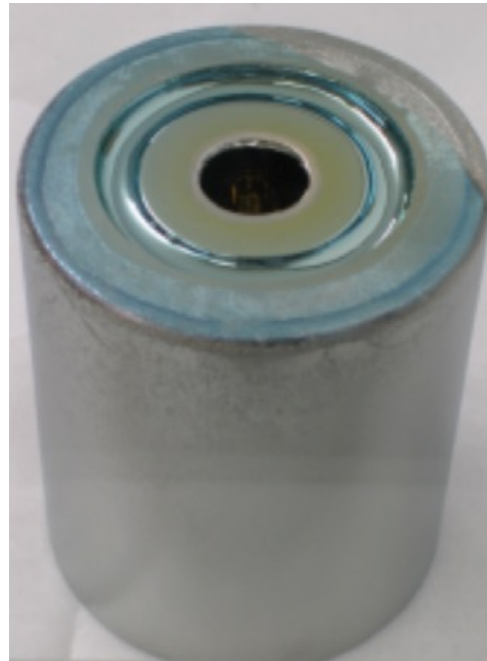
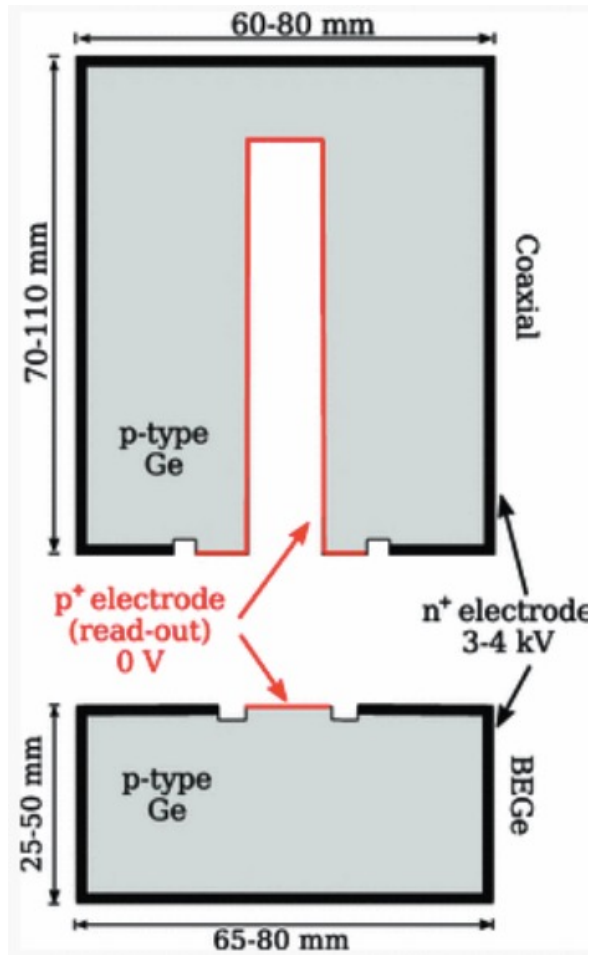
A. Bettini Padova University and INFN

Ge – diode

Technique well developed in gamma spectroscopy nuclear physics

Inverse polarisation, free charges produced by ionising radiation are collected and amplified

Need cryogenic temperatures



p⁺ contact = B implanted (thin)

n⁺ contact Li diffused

New in GERDA: naked in L Ar

Liquid Ar bath functions

- **Low temperature**
- **Passive shield**
- **Active veto for surface events**

GERDA. GE Detectors Array

Background sources and suppression

Strings of Ge diodes immersed in liquid Ar, acting as a cooler and a passive shield

The cryostat is immersed in ultrapure water tank shielding from environmental radiation

Background model supported by data

Close sources (supports, cables, electronics, L Ar contaminants)

^{214}Bi , ^{228}Th , ^{60}Co in the detector assembly

^{42}K from ^{42}Ar in L Ar bath

Residual ^{222}Rn dissolved in L Ar

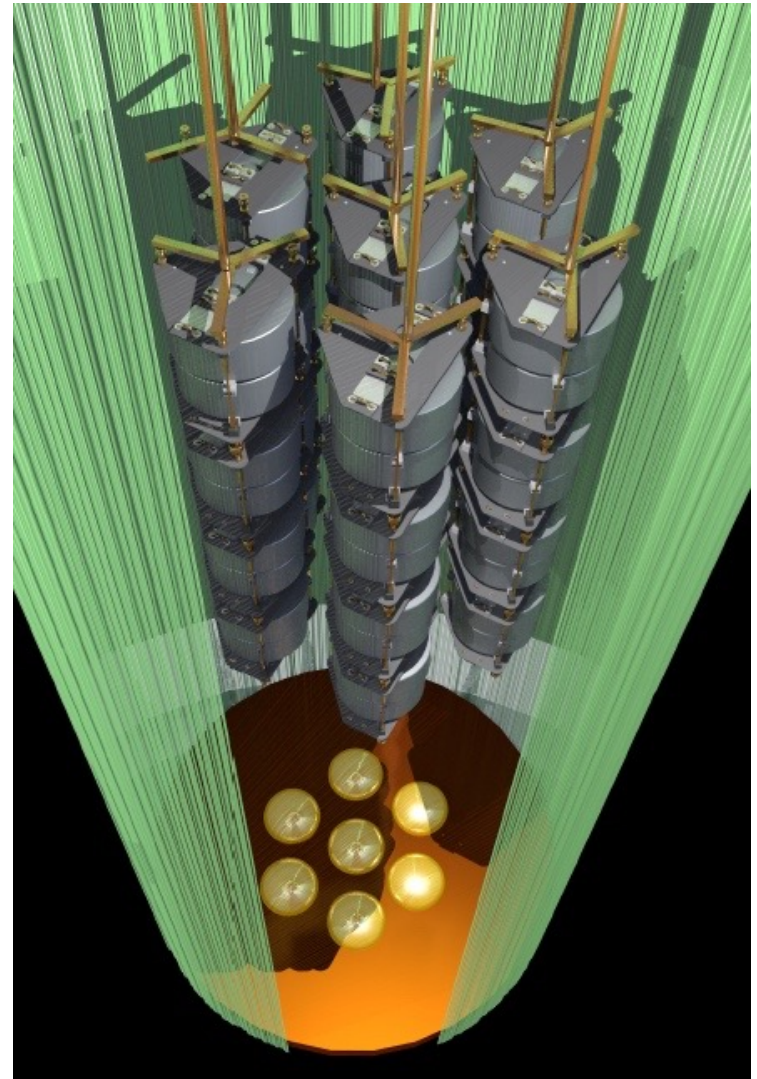
Surface α particles

Suppress by

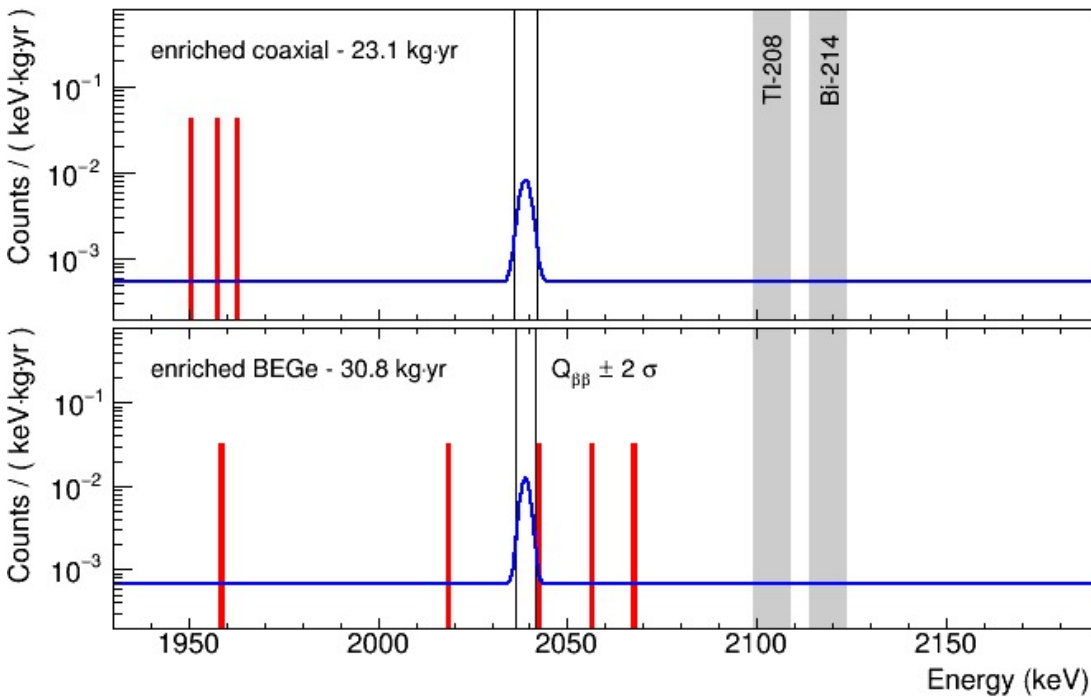
Anticoincidence between crystals

No light in the Ar (optical fibres)

Pulse shape analysis



GERDA. A background free experiment



M. Agostini et al PRL **125** 252502 (2020)

Exposure = 127.2 kg yr

$Q_{\beta\beta} = 2039.61 \pm 0.007$ keV

Energy resolution FWHM $\Delta E = 3.0 \pm 0.1$ keV (best)

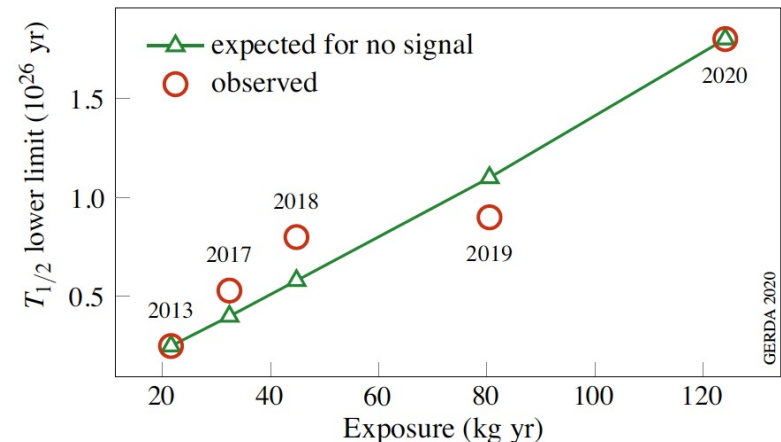
$b = 5.2 \times 10^{-4}$ cts/(keV kg yr)

$b \Delta E = 1.7$ counts/(FWHM t yr)

Median sensitivity $T_{1/2} > 1.8 \times 10^{26}$ yr (90% cl)

Limit $T_{1/2} > 1.8 \times 10^{26}$ yr (90% cl)

M_{ee} ^{76}Ge : <79-180 meV



LEGEND

Large Enriched Germanium Experiment for Neutrinoless $\beta\beta$ Decay (LEGEND). Members of GERDA, MAJORANA and others

Keep “**background free**” conditions while increasing Ge mass

200 kg. Starting (Infrastructures of GERDA)

Reduce background by 2.5

$b \Delta E = 0.7/(\text{FWHM t y})$

Electroformed copper (MD)

Increase light collection in L Ar

MD type F/E electronics

Larger crystals ICPC (2 kg)

$T_{1/2} > 10^{27}$ yr,

Exposure 1 t yr

1000 kg

$BI < 10^{-5}/(\text{keV kg y})$

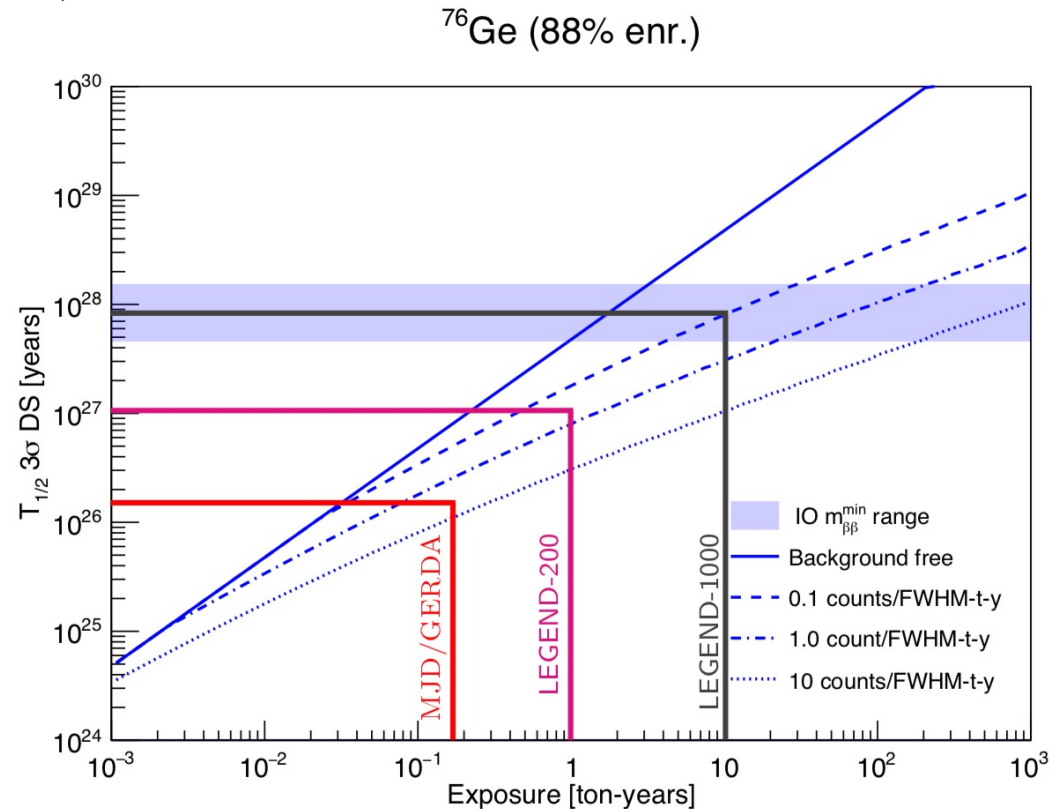
$b \Delta E = 0.025/(\text{FWHM t y})$

New infrastructure

Suppress ^{42}K from ^{42}Ar : depleted Ar

$T_{1/2} > 10^{28}$ yr

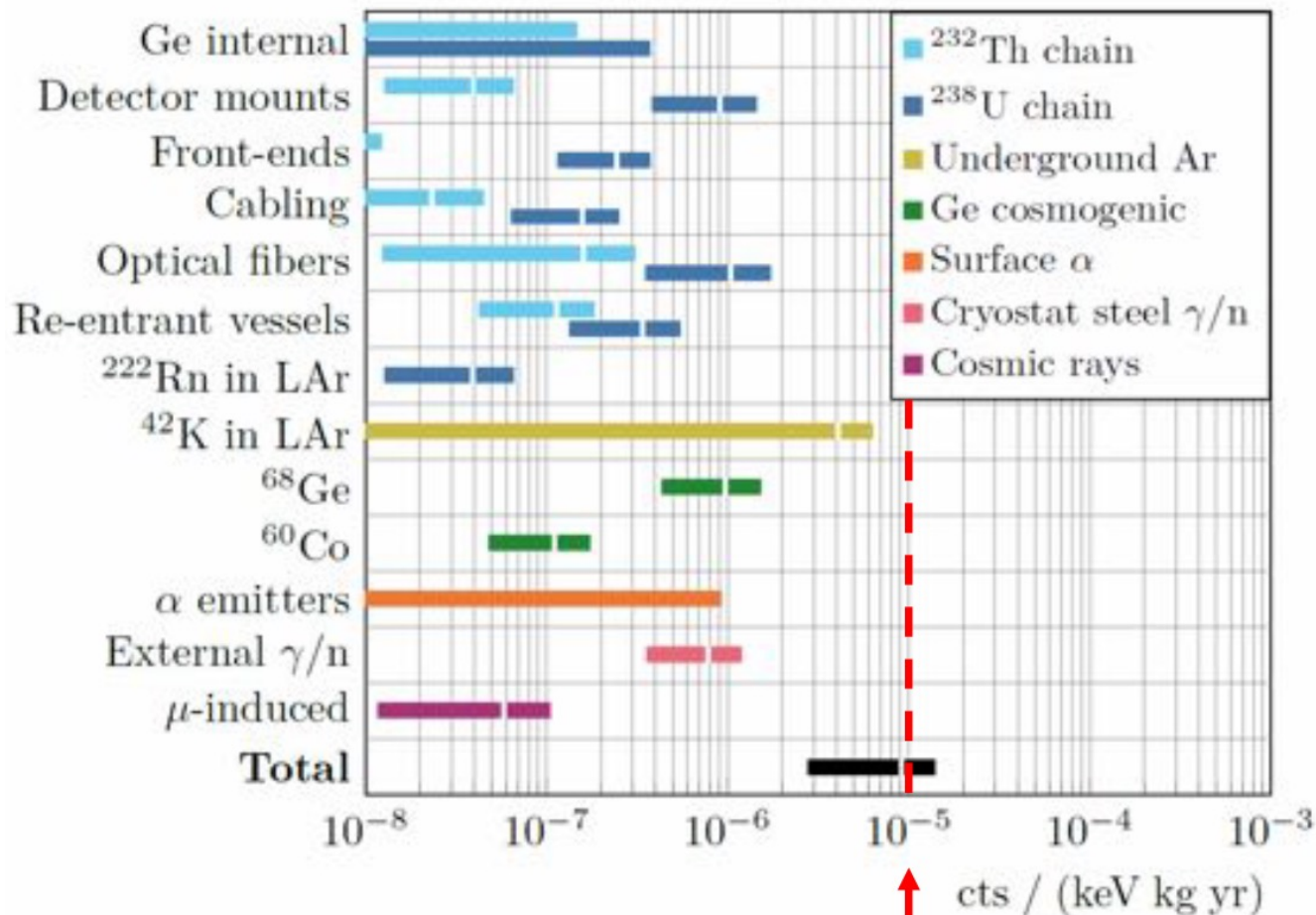
Exposure 10 t yr



arXiv:1810.00849

LEGEND 1000 Background model

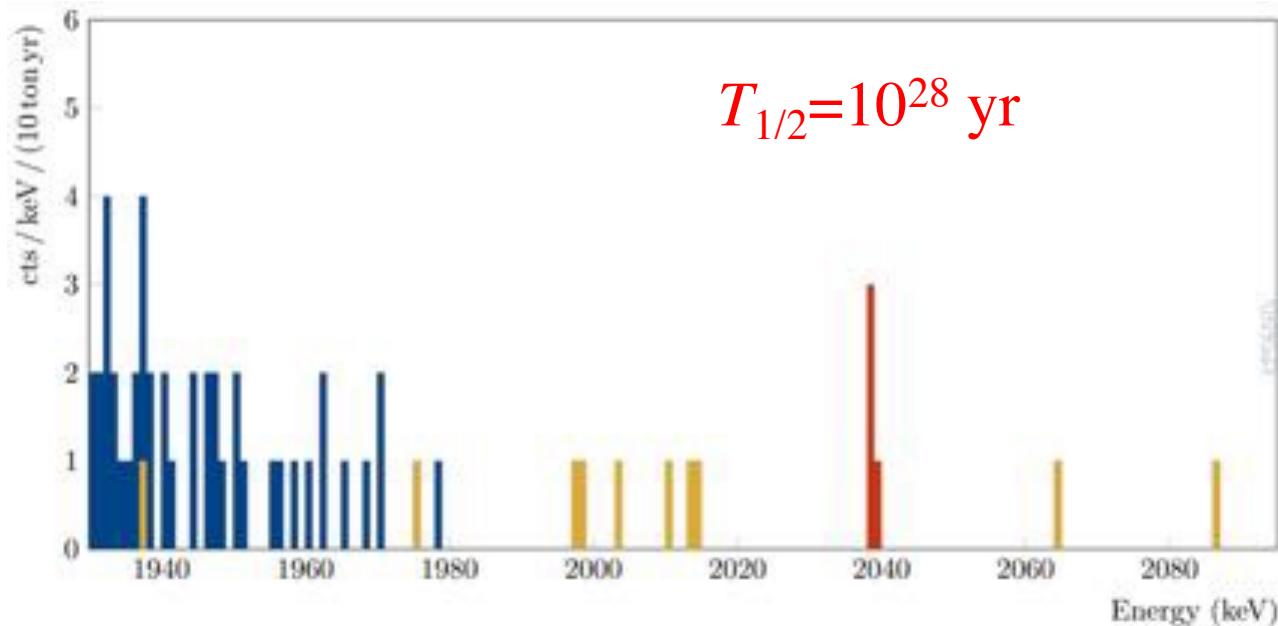
Background index at $Q_{\beta\beta}$ after all cuts



LEGEND-1000 background goal

LEGEND 1000 - $0\nu2\beta$ at 10^{28} yr

Simulated LEGEND-1000 spectrum after 10 yr exposure
with BI = 10^{-5} counts/(keV kg yr)



Flat residual background

No background peaks expected near $Q_{\beta\beta}$



KAMLAND-ZEN

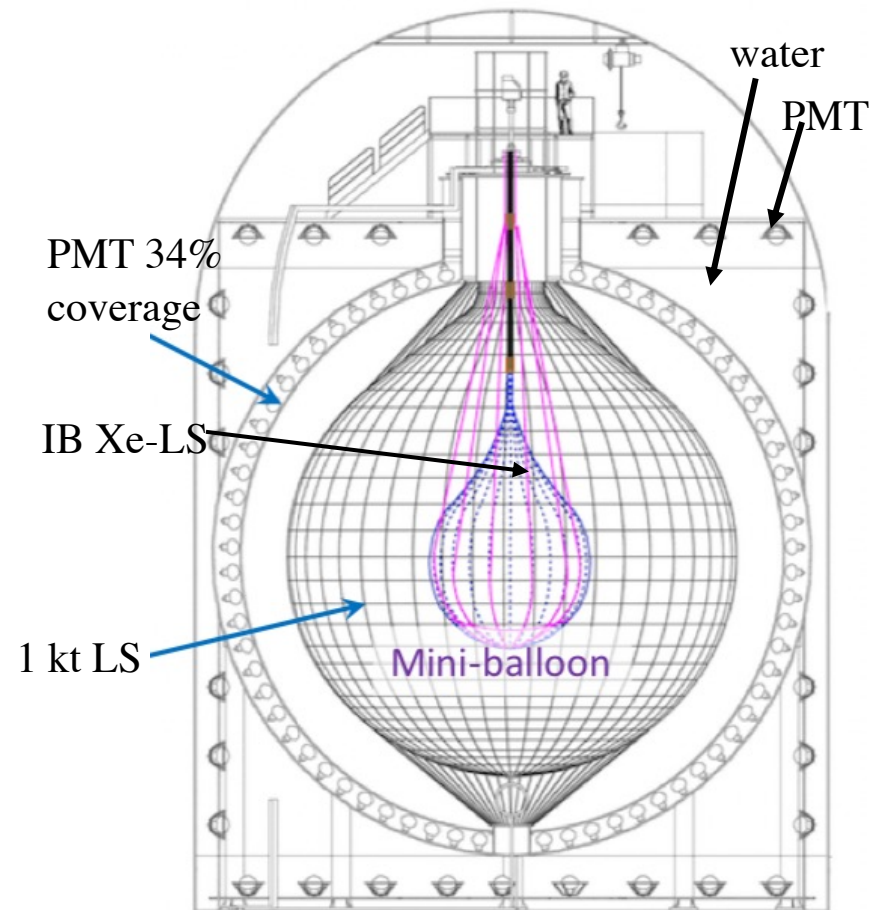
1994. R. Raghavan (1937-2011) proposal to load large masses of ^{136}Xe into a large-scale liquid scintillator detector PRL **72** 1411

Inner balloon contains ^{136}Xe doped LS
IB surface is source of background, reduced by geometrical cuts (**self shielding**)

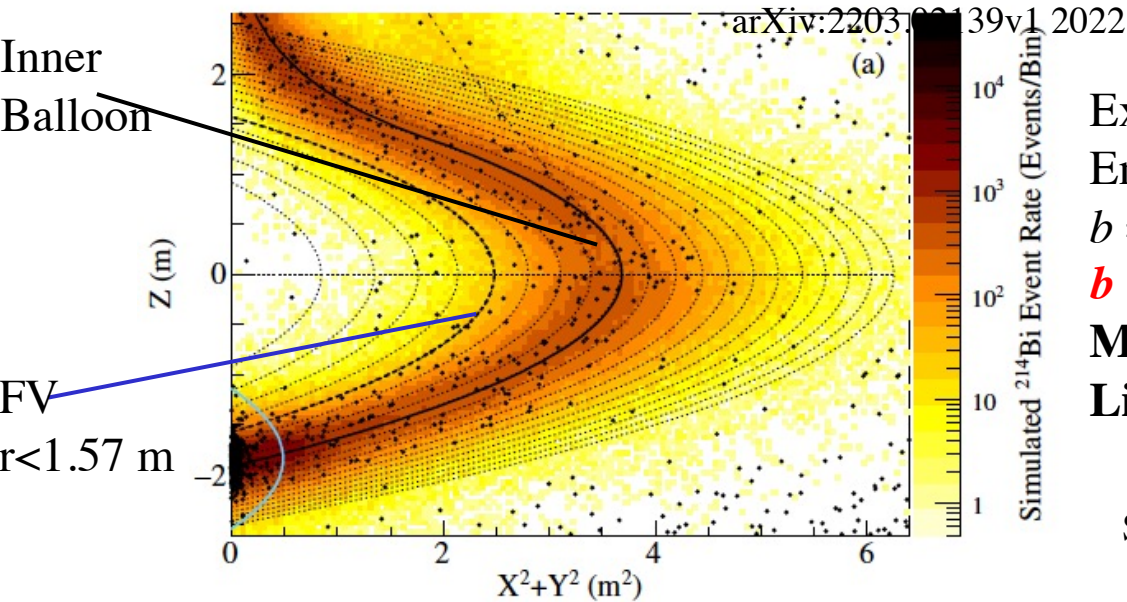
Data collection 5.2.2018-8.5.2021

^{136}Xe mass = 750 kg

Ultra clean IB

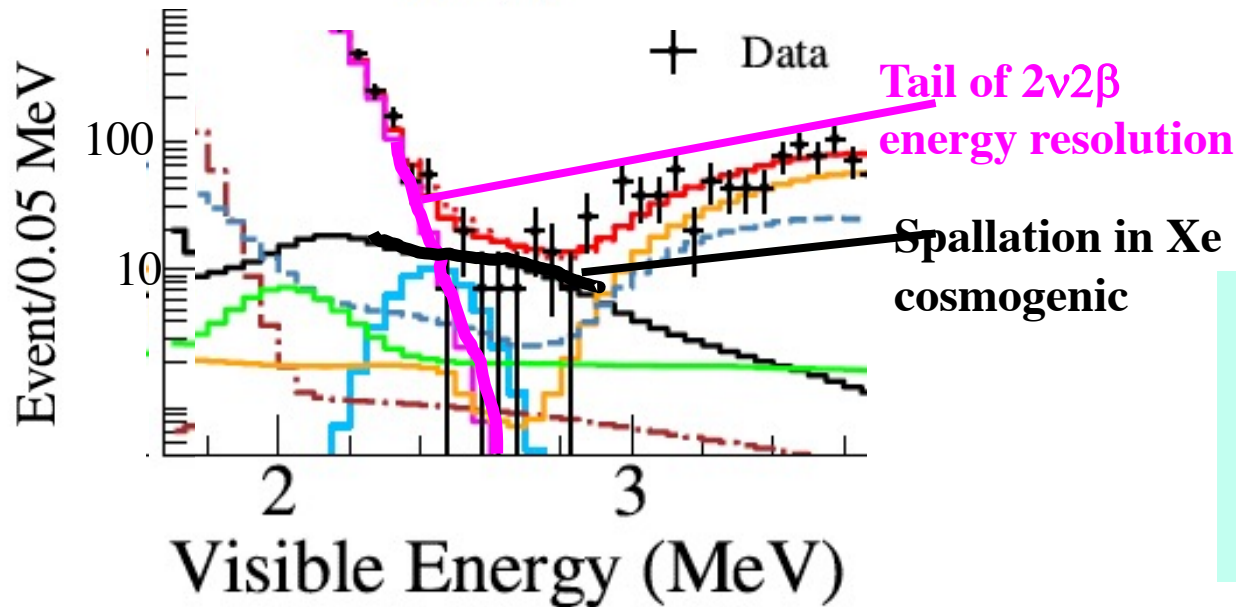


KAMLAND-ZEN-800



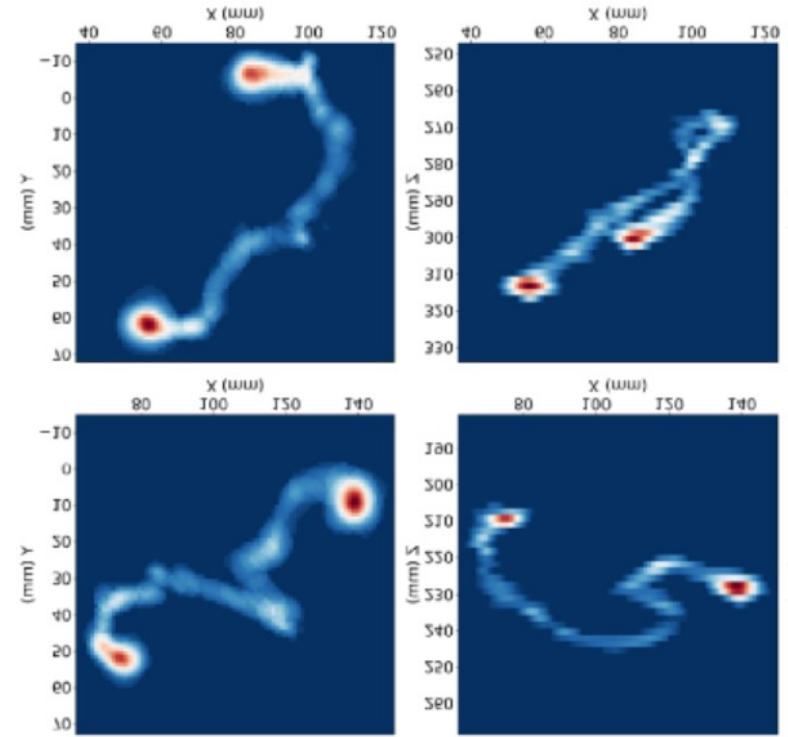
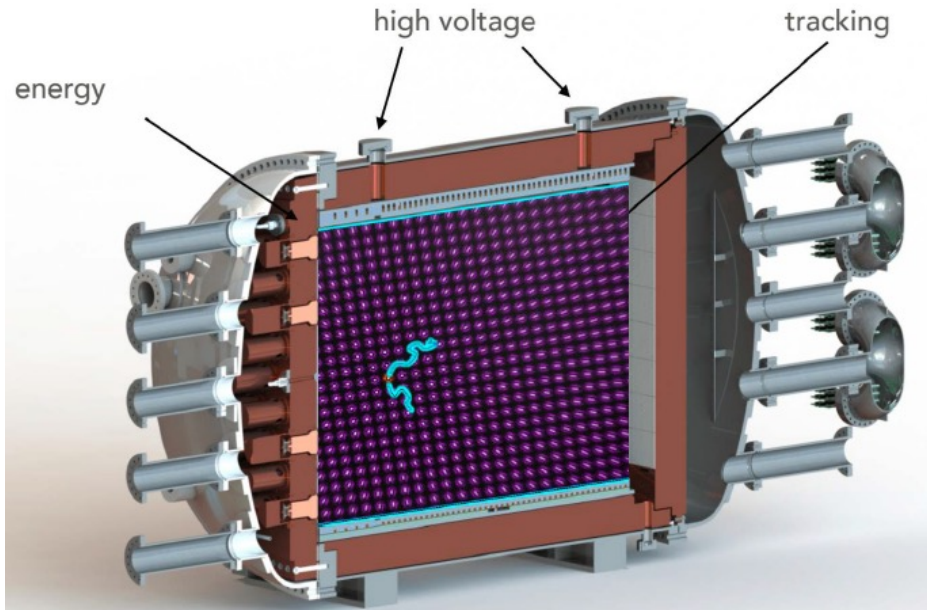
Exposure = 970 kg yr
 En. res FWHM $\Delta E = 265 \text{ keV}$ (poor)
 $b = 4 \times 10^{-5} \text{ counts/(keV kg yr)}$
 $b \Delta E = 11 / (\text{FWHM t yr})$
 Median sensitivity $T_{1/2} > 1.5 \times 10^{26} \text{ yr}$
 Limit $T_{1/2} > 2.3 \times 10^{26} \text{ yr}$ (90% cl)

S. Abe et al. arXiv:2203.02139

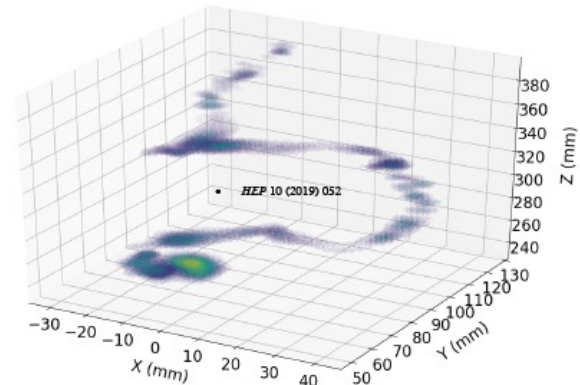


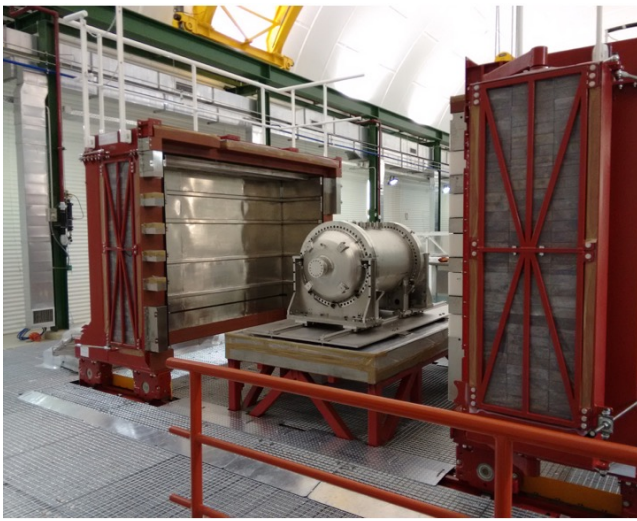
FUTURE: $T_{1/2} = 2 \times 10^{27} \text{ yr}$
 Improve light collection
 Reject film background
 (scintillating IB)
 Make entire volume FV

NEXT. Principle



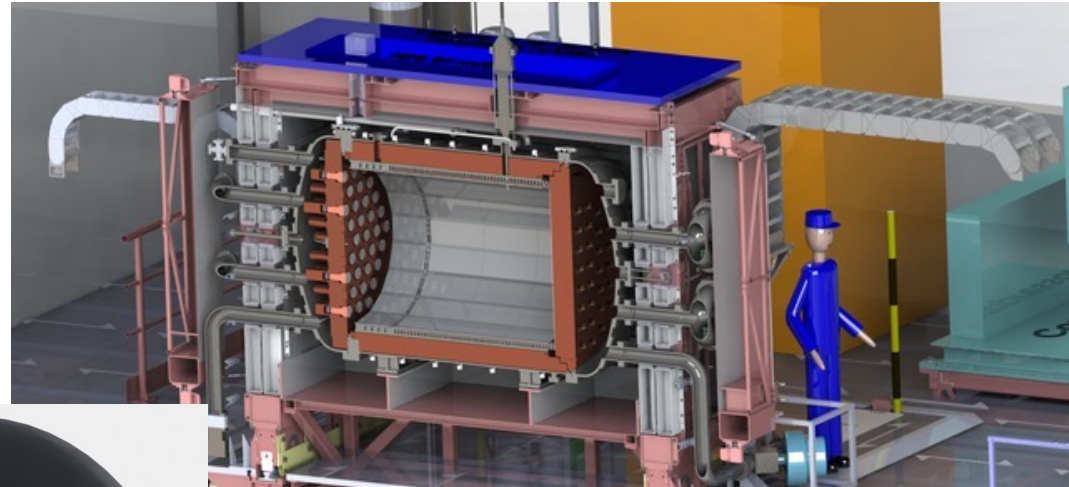
- High Pressure (10-15 Bar) gas Xe TPC
- Electro-Luminescence R/O
 $\Rightarrow \Delta E \text{ (FWHM)}/E < 1\%$
- Topological signature (30% efficiency)
- Strong potential for Ba tagging in situ
 $(^{136}\text{Xe} \rightarrow ^{136}\text{Ba}^{++} + 2 e^-)$





NEXT. The phases

NEXT-White. 10 kg ^{136}Xe : confirm design principles
 2018-2021 at LSC. runs with enriched and depleted Xe
 Energy resolution FWHM $\Delta E = 20 \text{ keV}$ (0.8%) (good)
 $b = 4 \times 10^{-4} \text{ counts}/(\text{keV kg yr})$
 JHEP 10 (2019) 052



NEXT-100.

Construction in the same infrastructure
 To run 2022-2026

$b \Delta E = 8 / (\text{FWHM t yr})$

Sensitivity $T_{1/2} \sim 10^{26} \text{ yr}$



NEXT-1t. 1000 kg TPC at 15 bar
 Central cathode plane and two R/O planes
 In water tank shield
Sensitivity with 5 t yr $T_{1/2} > 10^{27} \text{ yr}$

arXiv:2005.06467

y and INFN

Beyond 10^{27} . Look at the daughter



There's plenty of room at the bottom

Richard Feynman, 1960 (Nobel 1965)

There's even more room at the top

Jean-Marie Lehn, 1995 (Nobel 1987)



“Beyond molecular chemistry based on the covalent bond, there lies the field of **supramolecular** chemistry, whose goal it is to gain control over the intermolecular bond

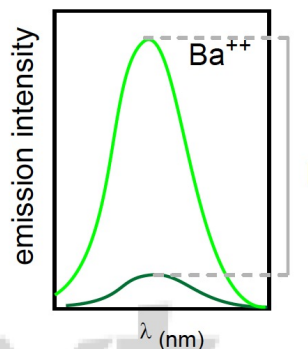
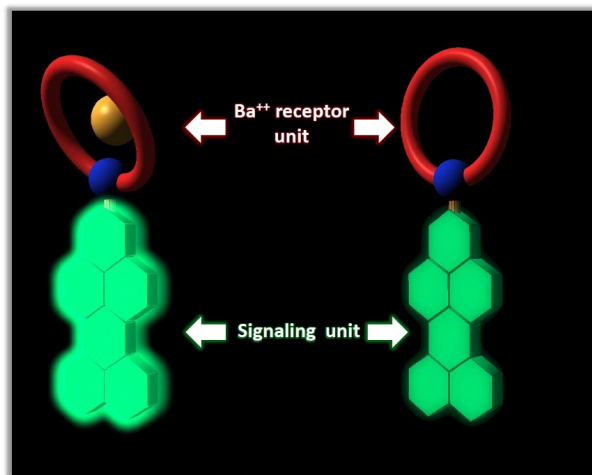
Build a supramolecular structure
able to **recognise** the Ba^{++}
daughter of the ^{136}Xe , **fix** it,
change shape to emit a
different light

NEXT-BOLD. Barium Observation Light Detector

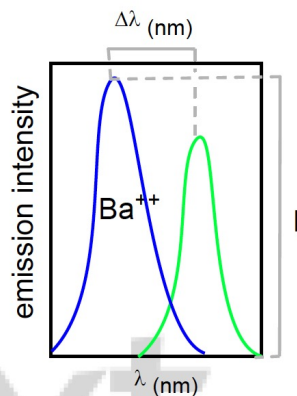
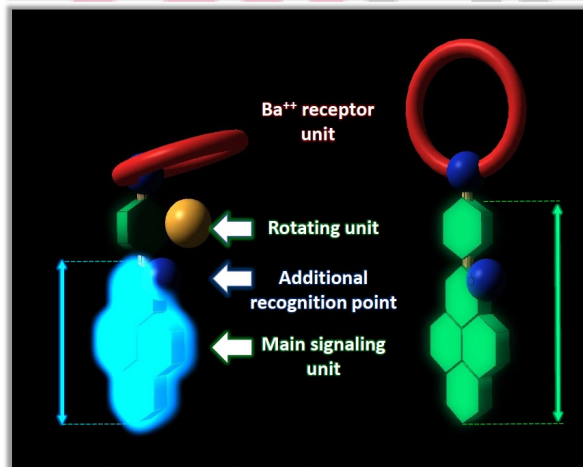
Exploit *single molecule fluorescent imaging* (SMFI) to visualize each single Ba ion as it reaches the TPC cathode

In ^{136}Xe gas phase (differently from in liquid phase) $\sim 100\%$ Ba^{++} remain ionised

Develop a supramolecular structure that change fluorescence light after chelating Ba dications
recognise Ba^{++} and **fix** (chelation) it (crown ether + fluorphore + linker}



For efficient detection, chelation should not simply induce a change of **intensity** of fluorescence light (upon laser illumination)



Adding a unit to rotate the Ba^{++} receptor, a change of the fluorescence light **wavelength** is obtained

<https://www.nature.com/articles/s41586-019-1169-4>

Fluorescence Bicolour Indicators

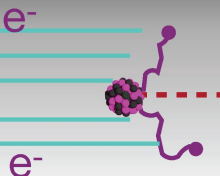
The FBI maintains its fluorescent properties in solid phase (silicagel)

Calculated discriminaton factor $\text{FBI} \cdot \text{Ba}^{++}$ vs $\text{FBI} = (2.5 \pm 0.6) \times 10^4$

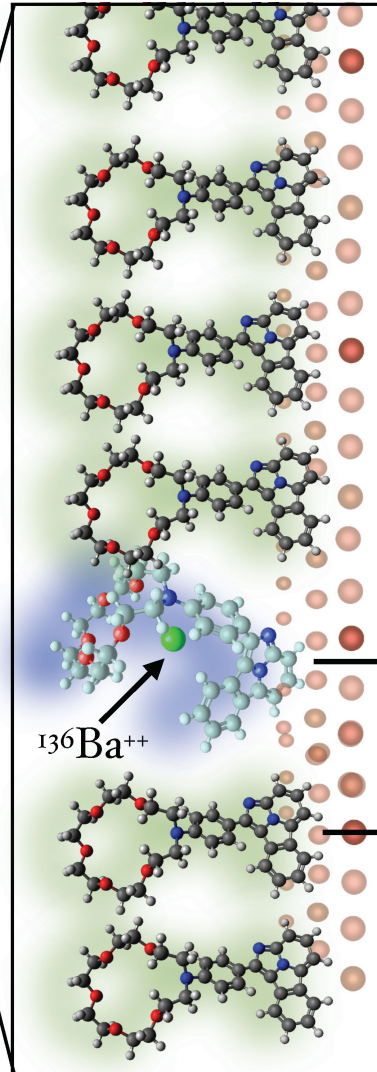
Barium tagging in NEXT:

A ML of Fluorescent Bicolor Indicators (FBI) traps the barium ion, thereby changing its optical properties.

Time Projection Chamber

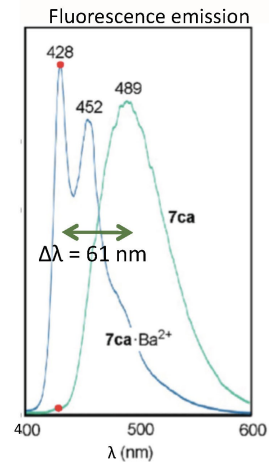


- First experiment of chelation in a dry medium:
- $\text{Ba}(\text{ClO}_4)_2$ was evaporated in UHV on silica pellets containing FBI. This changed its fluorescence from green to blue and proved that the trapping took place in UHV.

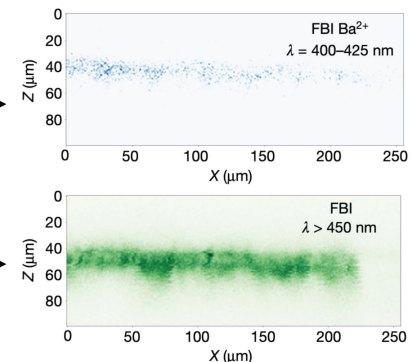
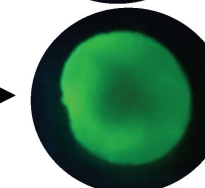
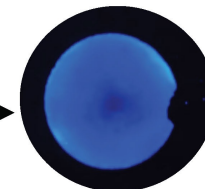


The Stokes shift between free FBI (cyan) and $\text{FBI} \cdot \text{Ba}^{2+}$ (blue) is larger than 35 nm.

This effect is specific of Ba^{2+} chelation and does not appear for other ions.



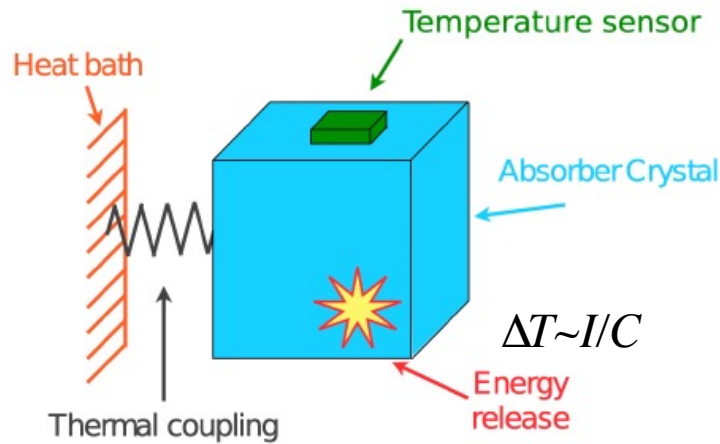
— FBI
— FBI-Ba



Microscopic tomographies were recorded for an unchelated (green) and a chelated (blue) pellet.

CUORE

^{130}Te in an array of 988 TeO_2 natural Te bolometers operating at about 10 mK at LNGS
Total mass = 741 kg (contributing to background). ^{130}Te mass about 206 kg (source of signal)

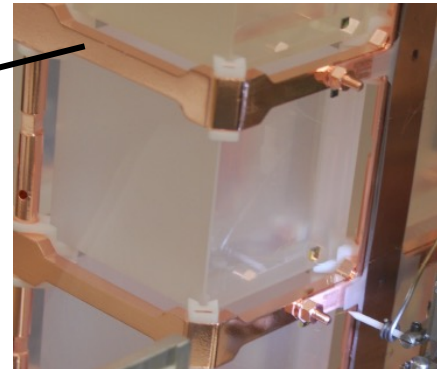


Voltage pulse induced in Ge thermistor used for R/O



$$b = (5.36 \pm 0.14) \times 10^{-2} / (\text{keV kg}_{\text{Te}130} \text{ yr})$$

Mainly degraded α from the surfaces of supports



2017-2021.

$$1038.8 \text{ kg}(\text{TeO}_2) \text{ yr} = 288.8 \text{ kg}(^{130}\text{Te}) \text{ yr}$$

En. Res. FWHM = 7.8 ± 0.5 keV (very good)

$$b \Delta E = 418 / (\text{FWHM } t_{\text{Te}130} \text{ yr})$$

Median sensitivity (90%): 2.8×10^{25} yr

Limit: $T_{1/2} > 2.2 \times 10^{25}$ yr

$M_{ee} < 90\text{--}305$ meV

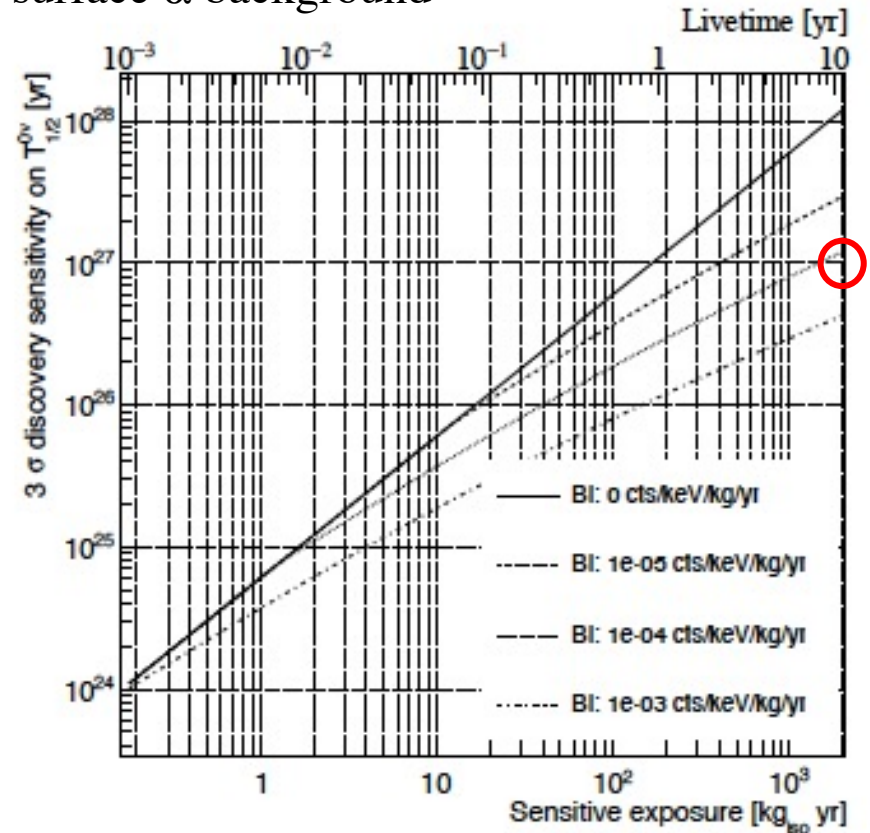
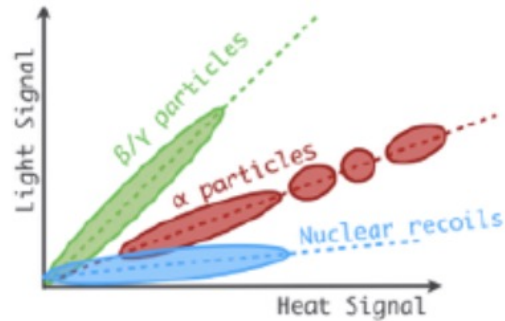
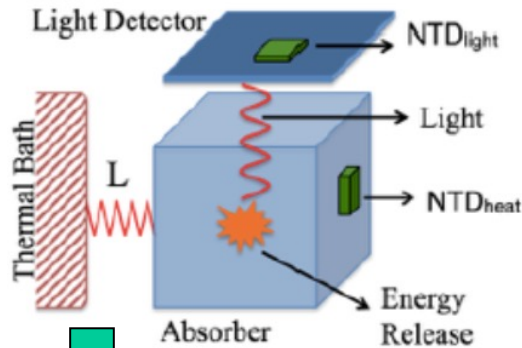
Will continue to 1000 kg yr ^{130}Te exposure

^{100}Mo CUPID

Will use the CUORE cryostat, after completion of the experiment

Source: $\text{Li}_2^{100}\text{MoO}_4$ scintillating crystals

Read out **HEAT and LIGHT**: identify and reject surface α background



Prototype experiments (CUPID-0, CUPID-Mo,...), R&D programme aiming at improve **b** by **1/100** and **ΔE** by **1/2** relative to CUORE

$$b \Delta E = 0.1 / (\text{keV } t_{100\text{MoO}_4} \text{ yr}) \times 4 \text{ keV} / (\text{FWHM } t \text{ yr})$$

Background free exposure of 2000 $\text{kg}(^{100}\text{MoO}_4) \text{ yr}$ (1120 $\text{kg}(^{100}\text{Mo})$) to reach 10^{27} yr sensitivity

The stage

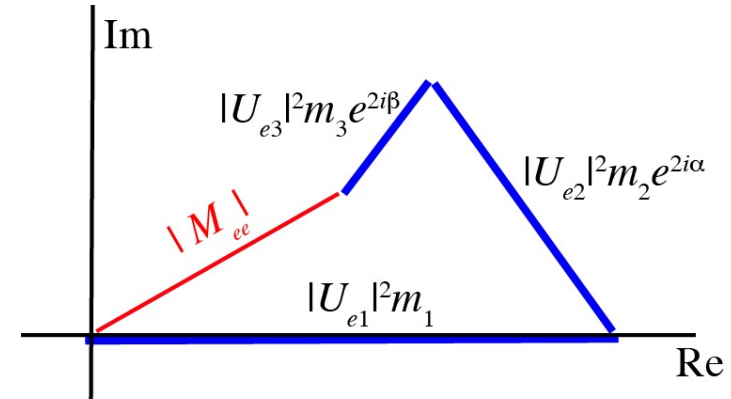
From cosmology $\Sigma m_i < 150\text{--}200$ meV model dependent

From $0\nu 2\beta$ limits on M_{ee}

^{136}Xe : $<36\text{--}156$ meV

^{76}Ge : $<79\text{--}180$ meV

^{130}Te : $<90\text{--}305$ meV



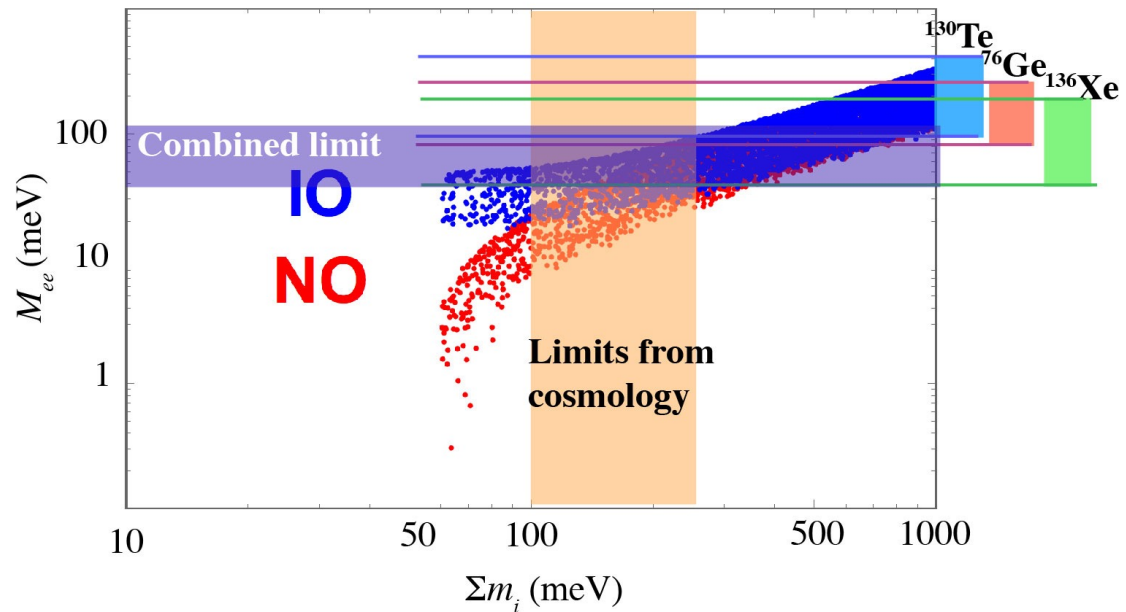
IO and NO bands correspond to Majorana phases $0\text{--}2\pi$

Dots from A. Ianni assuming random distribution of the phases
Combined 90% limit is from Lisi & Marrone arXiv:2204.0956

Depending on NME

$g_A = 1.276$

$$M_{ee} < 35.4 - 117.7$$



dreams



Look at the bottom



There's plenty of room at the bottom

of the energy scale

CNB: the cosmic sea of primordial **neutrinos** around us: **336/cm³**

$$T_\nu = (4/11)^{1/3} T_\gamma = 0.256 \text{ meV} = 1.95 \text{ K}$$

When neutrinos are non-relativistic the distinction between Dirac and Majorana becomes relevant
CNB, a laboratory to study neutrinos in unique kinematic regime

Dreaming on Cosmic Neutrino Background

In the hot and dense primordial universe neutrinos are in thermic equilibrium in the plasma

Having been produced by weak interactions, they were **flavour eigenstates** (6 in total) and **chirality eigenstates** with $\gamma_5 = -1$

Flavour states soon decohere into **the mass eigenstates**, which populate the present CNB

Neutrinos had almost only helicity $= -1$, antineutrinos $= +1$

When the temperature dropped below $\sim 1\text{MeV}$ neutrinos “froze out” from the plasma (CNB), **freely streaming** to us.

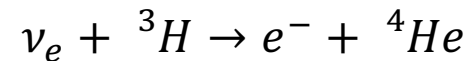
Chirality does not commute with the free Hamiltonian, it was not conserved

Helicity was conserved, being the particles non-interacting

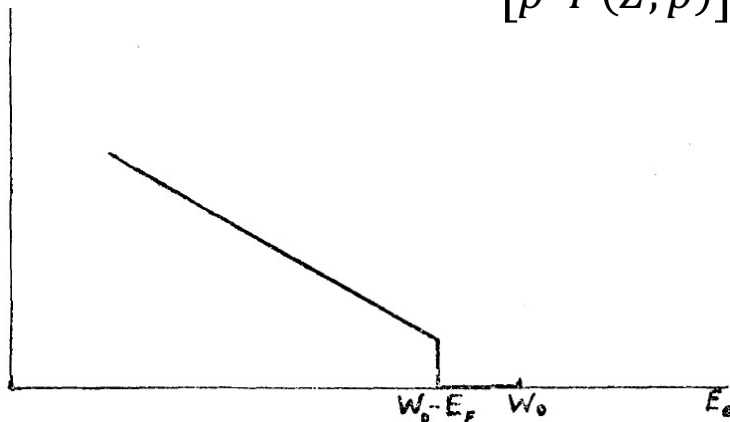
S Weinberg – Look at the end-point



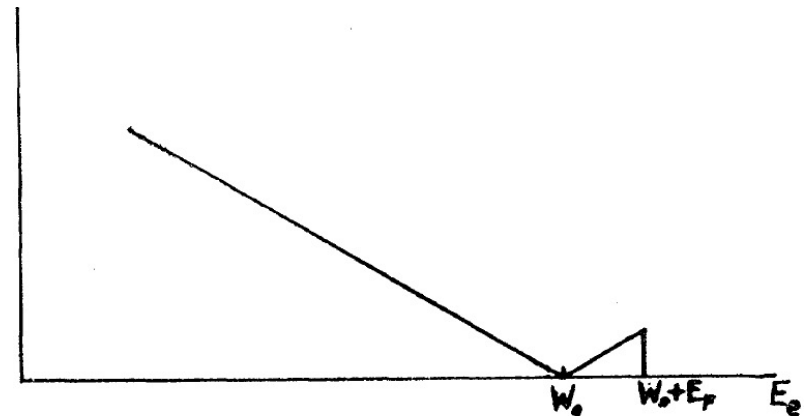
1962 S. Weinberg (1933 – 2021)
proposes CNB detection via inverse
beta decay, capture by a beta decaying
nucleus, since this is a **zero-threshold**
process *ibid.* Phys. Rev. **128** (1962) 1457



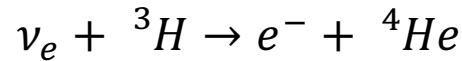
Kurie function $K = \left[\frac{N(p)}{p^2 F(Z, p)} \right]^{1/2}$



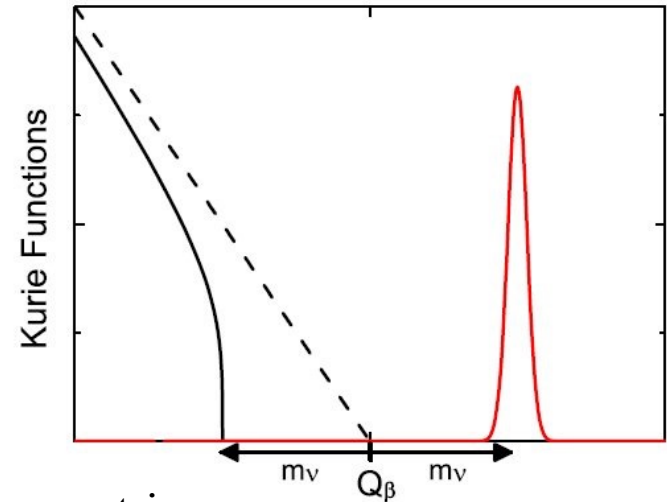
Electron neutrino mass KATRIN experiment



Detecting CNB



Consider one neutrino only
A peak at twice the neutrino mass
above the end-point of the beta-
decay spectrum should be observed



Capture rate is **twice** as large for Majorana than for Dirac neutrinos

CNB contains negative helicity neutrinos and positive helicity antineutrinos. Being non-relativistic, both helicity states contain a negative chirality component, hence both can interact.

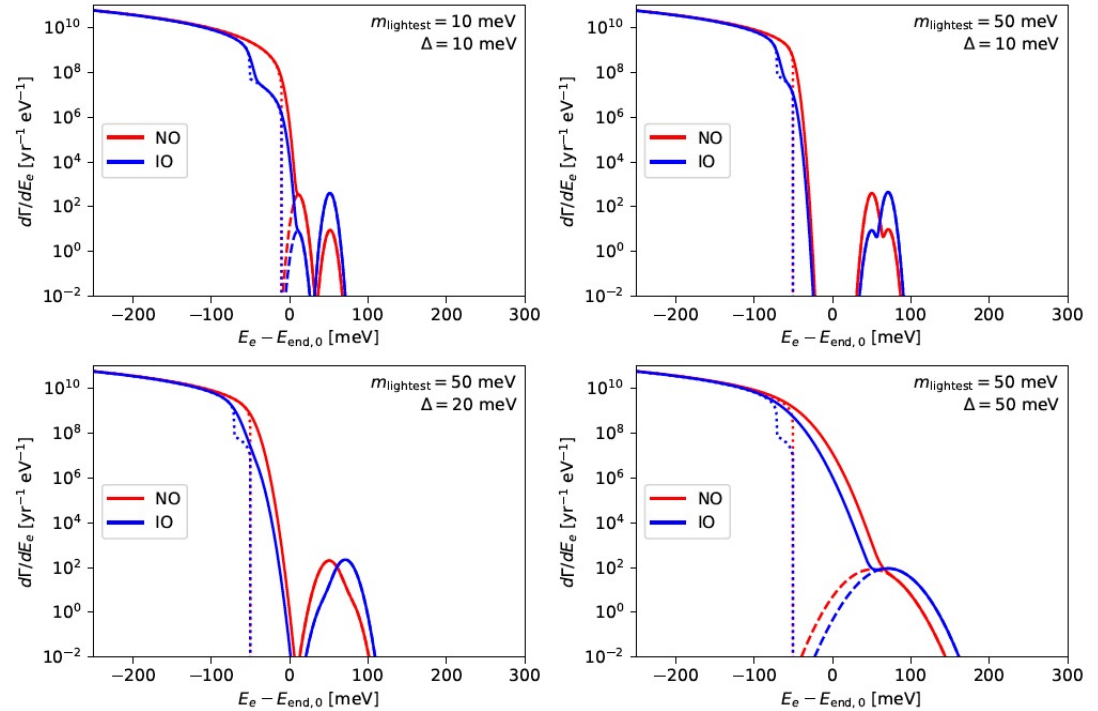
Dirac. The positive helicity antineutrino cannot be captured because $\bar{\nu}_e + p \rightarrow e^+ + n$ is energetically forbidden

Majorana. "Neutrino" and "antineutrino" are the same particle. Both helicity Majorana neutrinos interact weakly and can be captured

Challenges

Event rate. ${}^3\text{H}$ beta decay: about 10^{15} Bq/g. Target mass >100 g; presently impossible

Background from beta decay
Energy resolution should be
better than neutrino mass
 $\Delta E < m_i$



M. J. Betti et al. JCAP 07 (2017) 047 arXiv:1902.05508v1

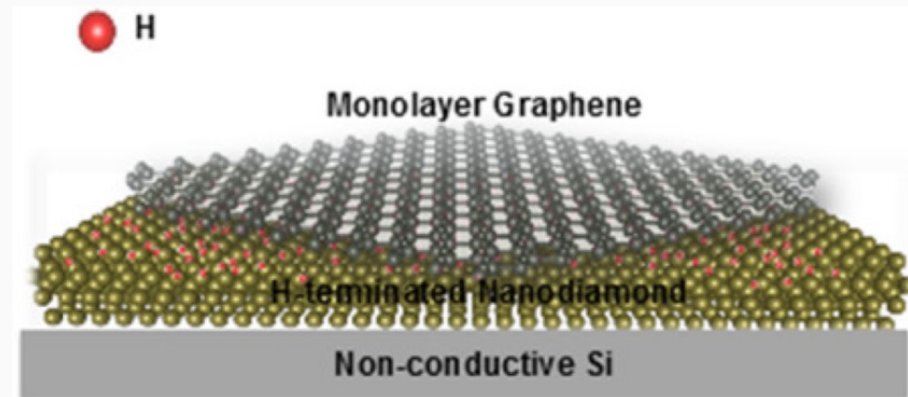
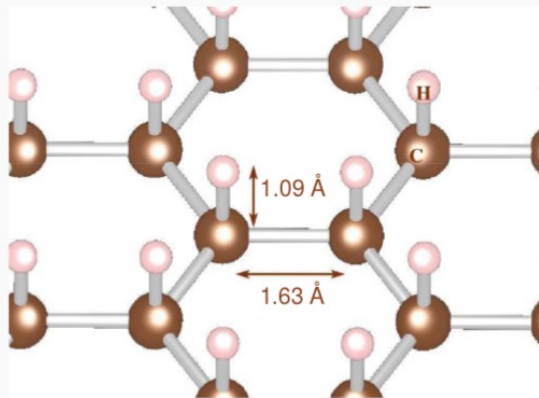
PTOLEMY @ LNGS

Be brave, dream the impossible
The PTOLEMY project
Prototypes forest being built



*There's even more room
at the top*

Target = supramolecular link tritium atoms on graphene monolayr



- Monolayer graphene consists of carbon atoms arranged in a two-dimensional honeycomb lattice;
- Graphene substrates are suitable to hold monoatomic Tritium layer through chemical absorption;

From Andrea Giachero

Low Energy Frontier RENP

Neutrino mass scale : few 10^{-2} eV

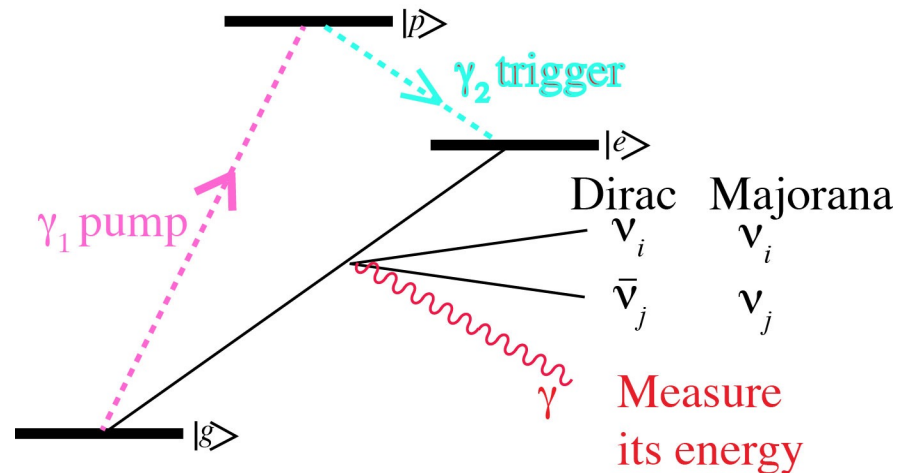
Atomic levels: **few eV** $\Rightarrow E/m$ order of 10^{-2} . Option for neutrino mass spectroscopy

Molecular electronic levels: **few 0.1 eV** $\Rightarrow E/m$ order of 10^{-1} . Majorana phases via interference?

Research programme at Okayama led by Yoshimura since some 20 years

RENP = Radiative Emission of Neutrino Pairs

QED forbidden by Parity



Rates extremely small



AVOGADRO squared is 10^{53}

Coherence on macroscopic dimensions

From N to N^2 MACROCOHERENCE

Rate amplification by macroscopic coherence

Dicke super-radiance (Phys Rev **93** (1954) 99-110): coherent volume is proportional to $L \lambda^2$

Coherent de-excitation develops from stimulated decays along pulse propagation direction

Observed: Skibowitz et al, PRL **30** (1973) 309

$$\text{Rate} \propto \left| \sum_j^N e^{i\vec{k} \cdot \vec{r}_j} \times M_{(e \rightarrow g)} \right|^2 \propto N^2 \quad (\text{for } |r_i - r_j| \leq \lambda)$$

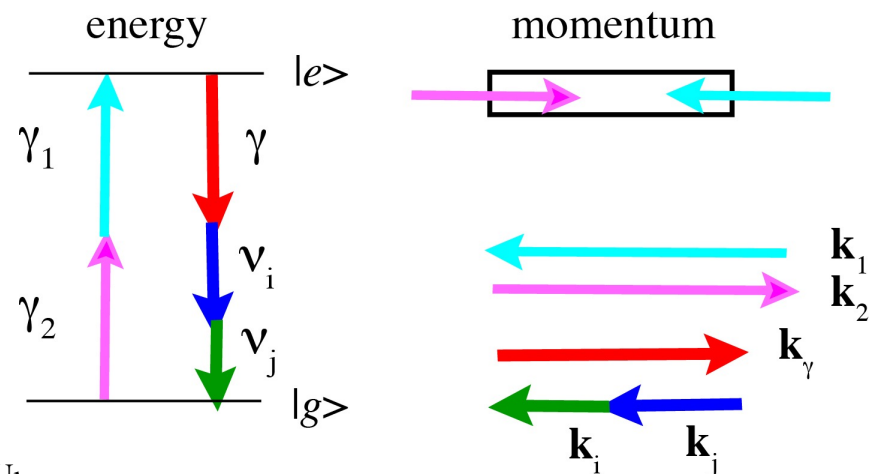
For plural outgoing/ingoing particles phase matching condition is obtained through momentum conservation

Coherent volume NOT limited by λ , can be macroscopic

$$\text{Rate} \propto \left| \sum_j^N e^{i(\vec{k}_1 + \vec{k}_2 + \vec{k}_3) \cdot \vec{r}_j} \times M_{(e \rightarrow g)} \right|^2 \propto N^2 \quad (\text{for } \vec{k}_1 + \vec{k}_2 + \vec{k}_3 = 0)$$

Employ 3 (or more) LASERs:
2 in opposite directions for excitation
1 to trigger the decay

Typical trigger bandwidth
 $\Delta E = 10^{-6}$ eV

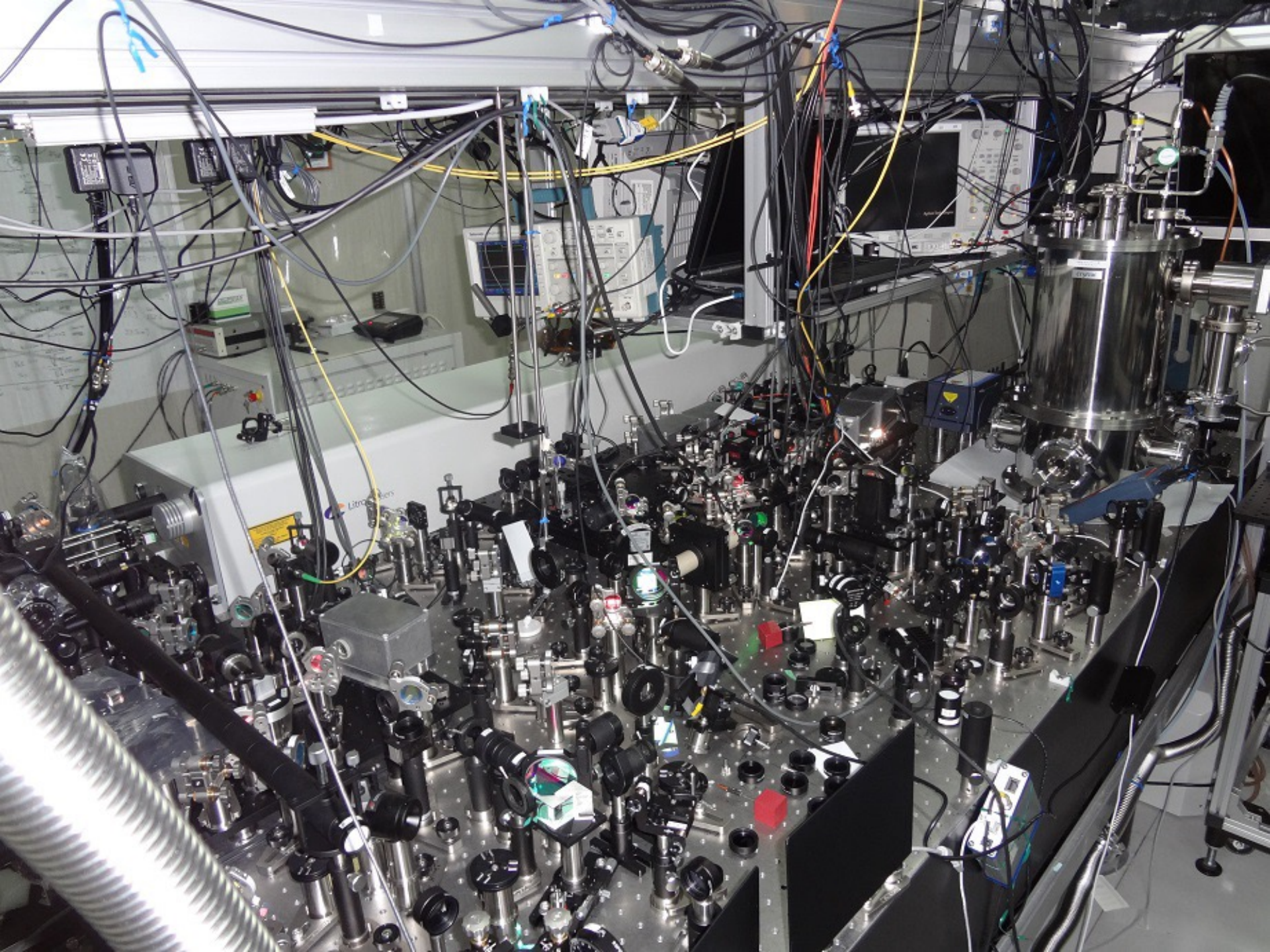


Conclusions

- The search for $0\nu 2\beta$ decay is still the main route to Majorana particles
- Present sensitivity in $T_{1/2}$ above 10^{26} yr (M_{ee} around 100 meV)
- 10^{27} yr (30 meV) in reach with ^{76}Ge , ^{136}Xe and ^{100}Mo in the current decennium (discovery possible)
- 10^{28} yr (10 meV) will need substantial further reduction of the background (discovery very likely)
- Additional research routes on the low energy frontier are being opened
 - Need deep understanding and development in atomic, molecular and supramolecular physics



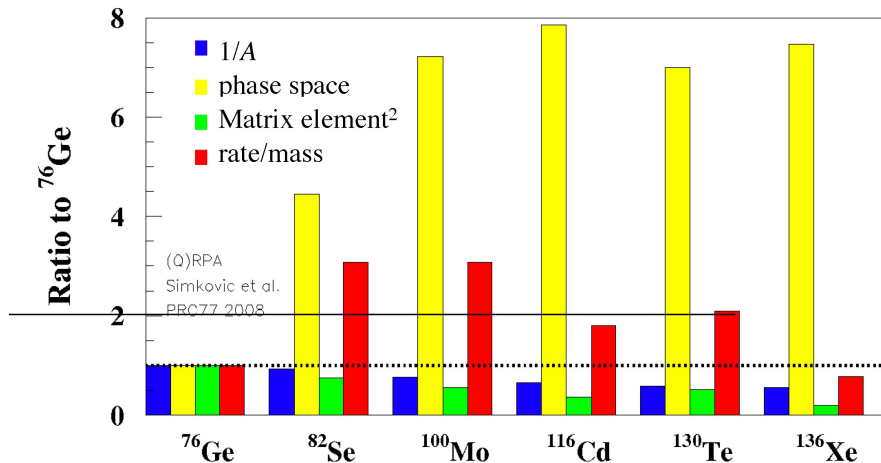
Thanks for your attention



Any favourable isotope?

Isotope	$Q_{\beta\beta}(\text{keV})$	Abundance (%)	$G (10^{-15} \text{ y}^{-1})$
^{76}Ge	2039.061(07)	7.8	2.36
^{100}Mo	3034.40(17)	9.6	15.91
^{130}Te	2527.5218(13)	33.8	14.20
^{136}Xe	2457.83(37)	8.9	14.56

PSF grows almost as Q^5 . Partially compensated by a decrease of NME with increasing A



No significantly favoured isotope
Within a factor 3
Choice of the isotope depends on
practical factors

PSF from F.F. Dappish et al. 2009.10119