

Model-independent study of vector-like quarks scenarios

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Erice

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- **Goal**: From a model-independent analysis, we want to obtain constraints on a general model featuring one or several VLQs.

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 - Presentation of VLQ
 - Phenomenology of VLQ
- 2 Visible decay
 - The XQCAT program
 - Study of the interferences
 - Offshellness analysis
- 3 Invisible decay (preliminary results)
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A *vector-like quark* (VLQ) is a quark whose left- and right-handed chiralities belong to the same representation of the symmetry group G of the underlying theory. For the Standard Model (SM), $G = SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$.

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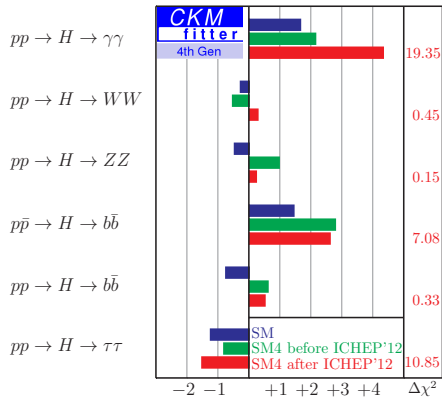
For VLQs, we have *both* left-handed and right-handed charged currents

$$J^{\mu+} = J_L^{\mu+} + J_R^{\mu+} = \bar{u}_L \gamma^\mu d_L + \bar{u}_R \gamma^\mu d_R = \bar{u} \gamma^\mu d = \textcolor{red}{V}$$

while for the SM chiral quarks we *only* have left-handed charged currents

$$J^{\mu+} = J_L^{\mu+} + J_R^{\mu+} \text{ with } \begin{cases} J_L^{\mu+} = \bar{u}_L \gamma^\mu d_L = \bar{u} \gamma^\mu (1 - \gamma^5) d = \textcolor{red}{V} - \textcolor{red}{A} \\ J_R^{\mu+} = 0 \end{cases}$$

Moreover, since a 4th generation of chiral quarks is excluded at 4.8σ by LHC Higgs data ([1209.1101]), searches for VLQs will acquire high priorities experimentally.



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Quantum numbers

	SM quarks $\begin{pmatrix} u \\ d \end{pmatrix} \begin{pmatrix} c \\ s \end{pmatrix} \begin{pmatrix} t \\ b \end{pmatrix}$	Singlets $(T)(B)$	Doublets $\begin{pmatrix} X \\ T \end{pmatrix} \begin{pmatrix} T \\ B \end{pmatrix} \begin{pmatrix} B \\ Y \end{pmatrix}$	Triplets $\begin{pmatrix} X \\ T \\ B \end{pmatrix} \begin{pmatrix} T \\ B \\ Y \end{pmatrix}$
$SU(2)_L$	$q_L = 2$ $q_R = 1$	1	2	3
Y	$Y_{q_L} = 1/6$ $Y_{u_R} = 2/3$ $Y_{d_R} = -1/3$	2/3 -1/3	7/6 1/6 -5/6	2/3 -1/3
$\mathcal{L}_m$	forbidden ¹	$-M\bar{\psi}\psi$	$-M\bar{\psi}\psi$	$-M\bar{\psi}\psi$

¹The Higgs mechanism is needed.

Visible and invisible decays

Two kinds of model:

- VLQ decaying to visible particles: the possibilities of decay for a VLQ T are $Z u_i$, $H u_i$ and $W^+ d_j$.



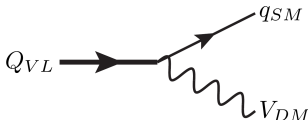
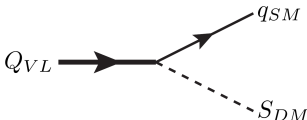
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- VLQ decaying to invisible particles (Dark Matter): the only possibility of decay for a VLQ T is χu_i , where χ is a DM particle (scalar or vector) made stable by a $\mathbb{Z}_2$ symmetry imposed to the Lagrangian.



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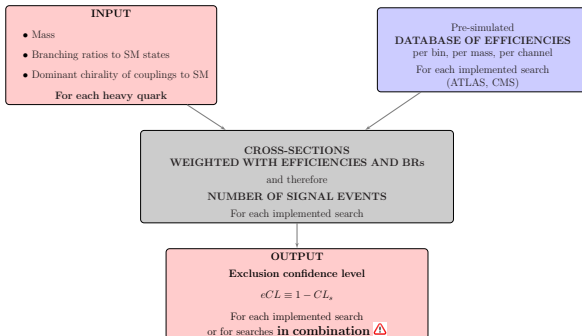
Presentation of the project

XQCAT in a nutshell

XQCAT = eXtra Quark Combined Analysis Tool

<https://launchpad.net/xqcat>

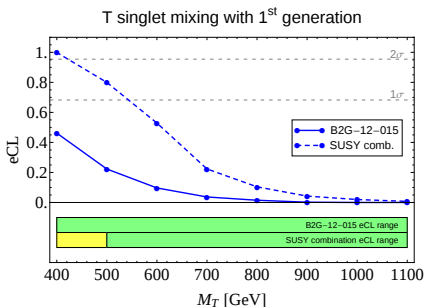
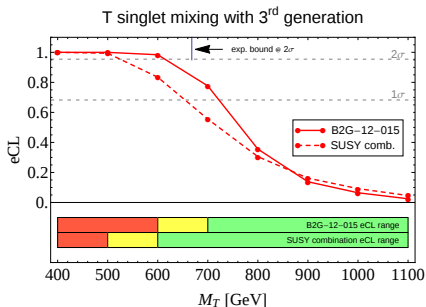
- 1) D. Barducci, A. Belyaev, M. Buchkremer, G. Cacciapaglia, A. Deandrea, S. De Curtis, J. Marrouche S. Moretti and LP, *Model Independent Framework for Analysis of Scenarios with Multiple Heavy Extra Quarks*, [arXiv:1405.0737 \[hep-ph\]](#) (submitted to JHEP)
- 2) D. Barducci, A. Belyaev, M. Buchkremer, J. Marrouche, S. Moretti and LP, *XQCAT: eXtra Quark Combined Analysis Tool*, [arXiv:1409.3116 \[hep-ph\]](#) (to be submitted to CPC)



First results of XQCAT: 1 T singlet

but with different mixing structure

$$BR(Zq) = BR(Hq) = 25\% \quad BR(Wq) = 50\%$$



- Stronger bounds when mixing with 3rd generation.
- Mixing with light generation: SUSY searches are more sensitive than direct searches.

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- Cross section: $\sigma \propto (\mathcal{A}_{T_1} + \mathcal{A}_{T_2})^2 \propto \underbrace{\mathcal{A}_{T_1}^2 + \mathcal{A}_{T_2}^2}_{\sigma_1 + \sigma_2} + \underbrace{2 \operatorname{Re}(\mathcal{A}_{T_1} \mathcal{A}_{T_2}^*)}_{\text{interference term}}$

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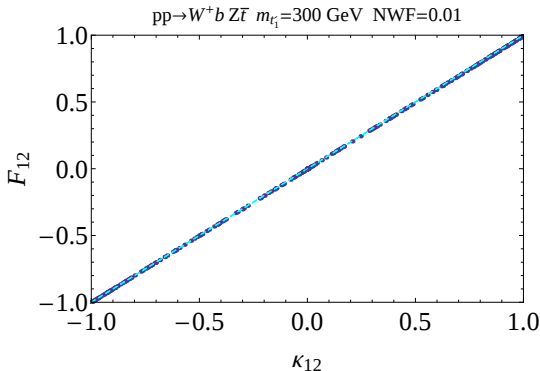
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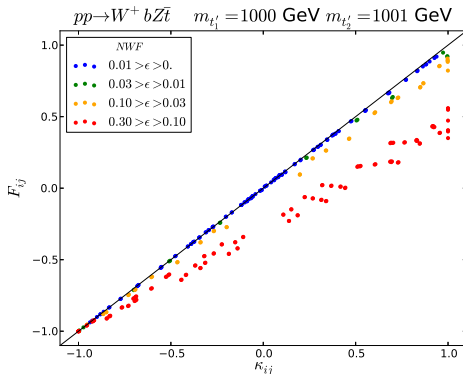
$$\Rightarrow \text{Ansatz: } \sigma_{\text{int}} \propto 2g_{1+}g_{1-}g_{2+}g_{2-} \operatorname{Re} \left\{ \left(\int \frac{dq^2}{2\pi} \mathcal{P}_1^0 \mathcal{P}_2^{0*} \right)^2 \right\}$$

Interference plot



$$F_{12} = \frac{\sigma_{\text{int}}}{\sigma_1 + \sigma_2} \simeq \frac{2 g_{1+} g_{1-} g_{2+} g_{2-} \operatorname{Re} \left\{ \left(\int \frac{dq^2}{2\pi} \mathcal{P}_1^0 \mathcal{P}_2^{0*} \right)^2 \right\}}{g_{1+}^2 g_{1-}^2 \left(\int \frac{dq^2}{2\pi} \mathcal{P}_1^0 \mathcal{P}_1^{0*} \right)^2 + g_{2+}^2 g_{2-}^2 \left(\int \frac{dq^2}{2\pi} \mathcal{P}_2^0 \mathcal{P}_2^{0*} \right)^2} = K_{12}$$

Limits of our ansatz

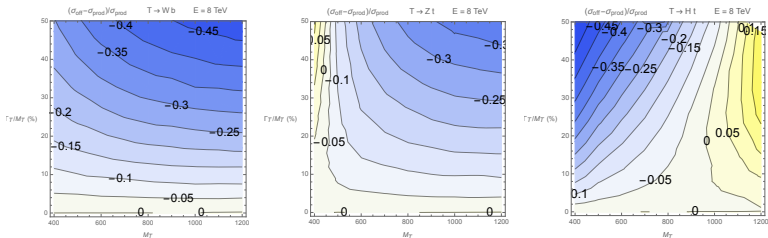


Interference effects cannot be well-treated for large $\Gamma/M \rightarrow$ let's explore in a quantitative way the deviations from NWA.

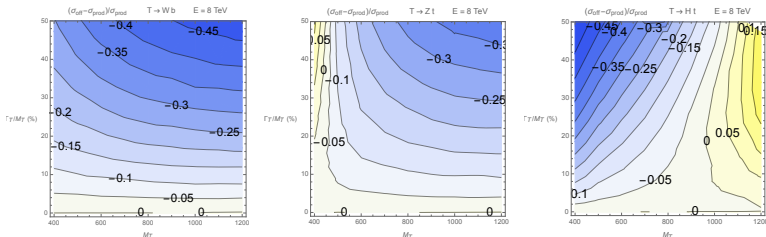
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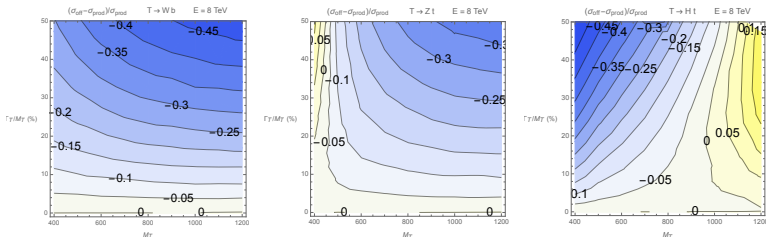


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- Wb and Zt channels: the offshell contribution is more important when Γ/M is large.
- Ht channel: different behaviour due to the fact that the coupling to the Higgs is proportional to the mass $M_T \rightarrow$ study of the differential distributions needed.

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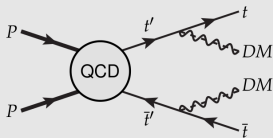
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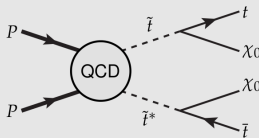
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- Interesting analogy with SUSY searches (same final state)

Decay into Dark Matter and top quarks



Heavy quark signal



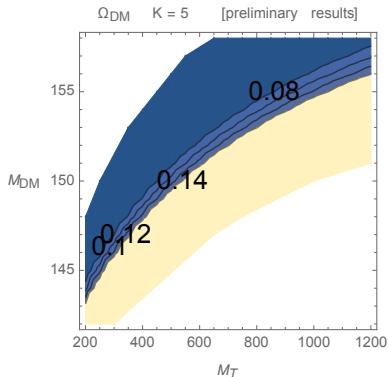
SUSY signal

If a signal is observed, it may be possible to distinguish a VLQ signal from a SUSY signal from the different kinematics of the events!

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By computing the relic density of DM in a model featuring a VLQ T decaying to a scalar DM particle we can impose strong constraints on the value of the masses and coupling.

$$\Omega_{DM} \sim \frac{1}{\langle \sigma v \rangle}$$



These preliminary results still have to be checked using micrOMEGAs.

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 - It may be possible to distinguish the origin of a signal by studying the *kinematics* of the event.

Thank you for your attention.

Backup slides

Dimension of representation

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$\Rightarrow \bar{q}HQ$ is a singlet

In term of representation of $SU(2)_L$ we have

$$① \quad q_R \otimes H \otimes Q = 1 \otimes 2 \otimes n = 1 \oplus \dots$$

or

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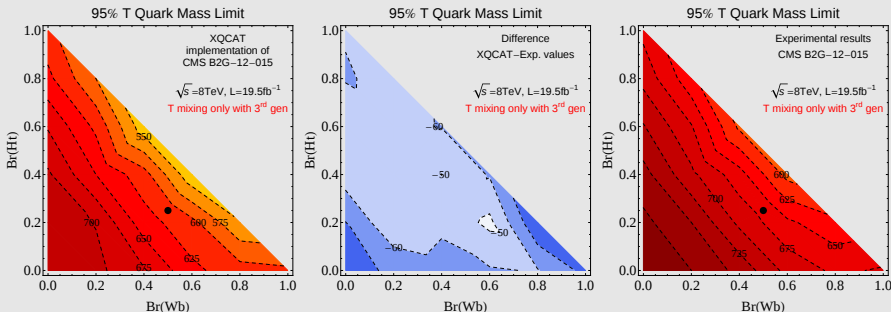
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Remark: We can also have higher representations for VLQs by considering model with more than one VLQ.

Validation of XQCAT: 1 T singlet

Validation plots: T mixing only with 3rd generation

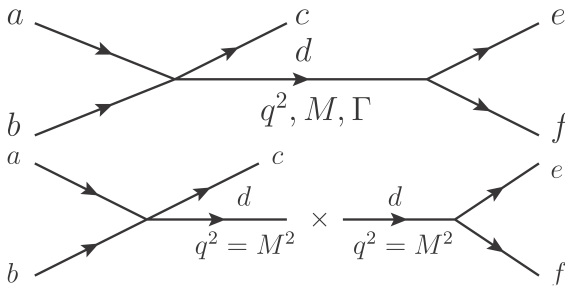


We reproduce CMS 95 % CL bounds within 50-60 GeV
in the whole BR range

Presentation of the NWA

The Narrow-Width Approximation (NWA) allows us to simplify the computation of complex processes $\rightarrow$ very useful and used in theoretical physics.

Basic idea: factorise the whole process into the on-shell production and the subsequent decay



Proof of the NWA

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With $\sigma = \frac{1}{F} \int d\Phi |\bar{\mathcal{M}}|^2$, $d\Phi = d\Phi_P \frac{dq^2}{2\pi} d\Phi_D$ and our approximation
 $\frac{1}{(q^2 - M^2)^2 + (M\Gamma)^2} \xrightarrow{\Gamma \ll M} \frac{\pi}{M\Gamma} \cdot \delta(q^2 - M^2)$, we find

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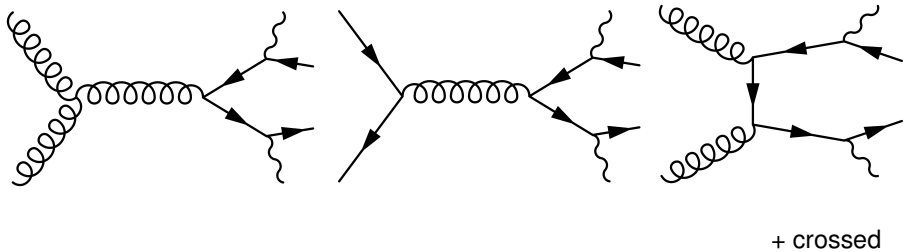
$$\mathcal{M} = \mathcal{M}_P \frac{1}{q^2 - M^2 - iM\Gamma} \mathcal{M}_D$$

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$$\sigma \simeq \sigma_P \cdot \frac{\Gamma_D}{\Gamma} \simeq \sigma_P \cdot BR$$

Generalisation



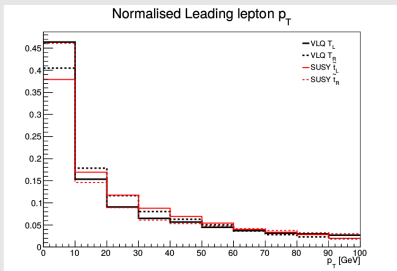
2 to 4 processes with fermionic propagators.

$$\sigma = \sigma_P \frac{\Gamma_{D^+}}{\Gamma_+} \frac{\Gamma_{D^-}}{\Gamma_-} = \sigma_P \cdot BR_+ \cdot BR_-$$

Comparison VLQ-SUSY [preliminary results]

Simulation of the process $pp \rightarrow t \bar{t} + \cancel{E_T}$ mediated by pair-produced T or $\tilde{t}$. Both signals processed through CheckMATE on a set of ATLAS searches with missing transverse energy in the final state.

One of the searches has a veto on leptons with p_T larger than 10 GeV



$$m_{t'/\tilde{t}} = 600 \text{ GeV}$$

$$m_{S_{DM}/\chi_0} = 10 \text{ GeV}$$

CheckMATE results: the SUSY point is excluded, but not the VLQ point!

The preliminary results looks promising.